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$L^{\infty}$ _error estimates for the obstacle problem revisited
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# $L^{\infty}$-ERROR ESTIMATES FOR THE OBSTACLE PROBLEM REVISITED 

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#### Abstract

In this paper, we present an alternative approach to a priori $L^{\infty}$-error estimates for the piecewise linear finite element approximation of the classical obstacle problem. Our approach is based on stability results for discretized obstacle problems and on error estimates for the finite element approximation of functions under pointwise inequality constraints. As an outcome, we obtain the same order of convergence proven in several works before. In contrast to prior results, our estimates can, for example, also be used to study the situation where the function space is discretized but the obstacle is not modified at all.


Key words. A priori error analysis, Linear finite elements, Obstacle problem

1. Introduction. This paper is concerned with a priori $L^{\infty}$-error estimates for the piecewise linear finite element approximation of the classical obstacle problem

$$
(P)\left\{\begin{array}{l}
\min \frac{1}{2} \int_{\Omega} \nabla v \cdot \nabla v \mathrm{~d} x-\langle f, v\rangle \\
\text { s.t. } v \in K:=\left\{z \in H_{0}^{1}(\Omega): z \geq \psi \text { a.e. in } \Omega\right\}
\end{array}\right.
$$

Pointwise error estimates for the problem $(P)$ have been studied by various authors before. They are typically derived by analyzing the error in mesh cells near the contact set of the continuous solution (i.e., the set where the solution and the obstacle coincide) and by subsequently applying the discrete maximum principle of Raviart-Ciarlet (cf. [5]). We only mention $[2,8,13,15]$ as references. In this paper, we take a more global perspective and demonstrate that a priori $L^{\infty}$-error estimates for the problem $(P)$ can also be obtained as corollaries of a more general stability result for discretized obstacle problems. Our method of proof has the advantage that the resulting error estimates are more flexible than their classical counterparts. They can, for example, also handle curved obstacles in the discrete setting. Moreover, our approach illustrates that the problem of estimating the approximation error for $(P)$ is, in fact, a problem of sensitivity analysis. This interpretation turns out to be very advantageous when it comes to analyzing the behavior of the approximation error in lower $L^{p}$-norms and the limitations of the piecewise linear finite element method. The alternative viewpoint provided by our analysis and the flexibility of our estimates were, for example, of major importance for the construction of two counterexamples found in a companion paper [4] which demonstrate that the convergence rates obtained for the $L^{\infty}$-error are - at least in the one-dimensional setting - also optimal if the $L^{p}$-error, $p>1$, is considered. The latter implies in particular that the Aubin-Nitsche trick does not work for the obstacle problem. We refer to [4] for details on this topic.
The outline of this paper is the following: In Section 2, we clarify the notation, address the used discretization scheme, and recall basic results about the solvability of the obstacle problem and the regularity of its solution. In the subsequent section, we introduce the concept of discrete supersolutions and use it to study the stability of the approximate problems obtained from the finite element discretization. We will see

[^0]here that, in contrast to the continuous setting, the solution operator of a discretized obstacle problem is not Lipschitz as a function of the obstacle if the latter is allowed to be curved. Section 4 is devoted to the a priori error analysis in the $L^{\infty}{ }_{-}$norm. Here, it is demonstrated that $L^{\infty}$-error estimates for the obstacle problem follow straightforwardly from the stability results of Section 3 if the Ritz projection of the continuous solution is identified with the solution of an appropriately defined discrete problem. The order of convergence that we ultimately obtain in this section is the same as in the classical works of Nitsche [15] and Baiocchi [2]. Lastly, in Section 5 we conclude our investigation with some remarks and a discussion of open problems. The appendix of this paper contains results about one-sided finite element approximations that are needed for our argumentation. The theorems found there may also be of independent interest.
2. Preliminaries. In what follows, $\Omega$ will always denote a bounded Lipschitz domain in $\mathbb{R}^{d}$, where $d \in \mathbb{N}$ is arbitrary but fixed. Furthermore, we will use the standard abbreviations $H_{0}^{1}(\Omega), W^{m, p}(\Omega), C^{m, \gamma}(\bar{\Omega})$ and $H^{-1}(\Omega)$ for the Sobolev spaces on $\Omega$, the Hölder spaces on the closure $\bar{\Omega}$ and the dual of $H_{0}^{1}(\Omega)$ w.r.t. the $L^{2}$-inner product. The pairing between elements of $H_{0}^{1}(\Omega)$ and $H^{-1}(\Omega)$ will be denoted with $\langle.,$.$\rangle . We refer to [1] and [7] for details.$
As already mentioned, the objective of this paper is to study the classical unilateral obstacle problem with zero boundary conditions: Given an $f \in H^{-1}(\Omega)$ (the force) and a measurable function $\psi: \Omega \rightarrow \mathbb{R}$ (the obstacle) find the solution to
\[

(P)\left\{$$
\begin{array}{l}
\min \frac{1}{2} a(v, v)-\langle f, v\rangle \\
\text { s.t. } v \in K:=\left\{z \in H_{0}^{1}(\Omega): z \geq \psi \text { a.e. in } \Omega\right\}
\end{array}
$$\right.
\]

The bilinear form $a$ appearing here is defined to be

$$
a: H_{0}^{1}(\Omega) \times H_{0}^{1}(\Omega) \rightarrow \mathbb{R}, \quad(v, w) \mapsto \int_{\Omega} \nabla v \cdot \nabla w \mathrm{~d} x
$$

Using that $a$ is coercive (due to the inequality of Poincaré-Friedrichs), it is easy to prove that $(P)$ admits a unique solution provided the admissible set $K$ is not empty. A detailed analysis shows the following:
THEOREM 2.1. If $\psi: \Omega \rightarrow \mathbb{R}$ is a measurable function such that $K$ is not empty, then for all $f \in H^{-1}(\Omega)$ there is a unique solution $u \in K$ to the problem $(P)$ and this solution is also uniquely determined by the variational inequality

$$
\begin{equation*}
u \in K: \quad a(u, u-v) \leq\langle f, u-v\rangle \quad \forall v \in K \tag{2.1}
\end{equation*}
$$

Moreover, the solution map $S: H^{-1}(\Omega) \ni f \mapsto u \in H_{0}^{1}(\Omega)$ is Lipschitz continuous. If, further, there exists a $2 \leq q<\infty$ such that $\psi, f$ and $\Omega$ satisfy

- $f \in L^{q}(\Omega), \psi \in W^{2, q}(\Omega)$ and $\operatorname{tr} \psi \leq 0$ a.e. on $\partial \Omega$,
- there exists a constant $C=C(\Omega, q)>0$ such that for all functions $v \in H_{0}^{1}(\Omega)$ with $\Delta v \in L^{q}(\Omega)$ it holds

$$
\begin{equation*}
\|v\|_{W^{2, q}} \leq C\|\Delta v\|_{L^{q}} \tag{2.2}
\end{equation*}
$$

then the solution $u$ is in $W^{2, q}(\Omega)$ and there exists a constant $C^{\prime}=C^{\prime}(\Omega, q)$ such that

$$
\|u\|_{W^{2, q}} \leq C^{\prime}\left(\|f\|_{L^{q}}+\|\max (-\Delta \psi-f, 0)\|_{L^{q}}\right)
$$

Proof. The unique solvability of $(P)$, the characterization of the solution $u$ by (2.1) and the Lipschitz continuity of the solution operator follow from standard results like the well-known theorem of Lions-Stampacchia (see, e.g., [12, Chapter II]). The $W^{2, q_{-}}$ regularity of the solution can be obtained with an approximation argument. We refer to [12, Chapter IV] for details.

A constant $C(\Omega, q)$ with property (2.2) exists, for example, if the domain $\Omega$ has a $C^{1,1}{ }^{-}$ boundary and $2 \leq q<\infty$ (see [9, Theorem 9.15 , Lemma 9.17 ]) or if $\Omega$ is a polygon with largest interior angle $\alpha$ and $2 \leq q<(1-\pi /(2 \alpha))^{-1}$ (see [10, Theorem 4.4.4.13] and [11, Theorem 2.2.3, Theorem 2.4.3.]). It should be noted that the solution $u$ to the problem $(P)$ will in general not possess higher derivatives than stated in the last theorem even if the obstacle $\psi$ and the force $f$ are smooth. If we consider, for example, the situation $\Omega=(-2,2), f(x)=0$ and $\psi(x)=1-x^{2}$, then the solution $u$ is a spline whose second derivatives are discontinuous at the boundary of the set $\{u=\psi\}$ where the solution and the obstacle coincide. This illustrates that higher order finite elements provide little practical advantages in the case of problem $(P)$ (at least as far as non-adaptive methods are concerned) and explains, why it makes sense to restrict the analysis to piecewise linear functions.

Having dealt with the existence, the uniqueness and the regularity of the exact solution, we now turn our attention to the discretization. First, let us recall some basic concepts (cf. [3]):

Definition 2.2. If $\Omega \subset \mathbb{R}^{d}$ is a bounded domain with a Lipschitz boundary, then a collection $\mathcal{T}=\left\{T_{i}\right\}$ of finitely many closed d-dimensional simplices $T_{i}$ is called $a$ triangulation of $\Omega$ if the following holds:

- $\bigcup T_{i}=\bar{\Omega}$,
- If $\mathcal{C}_{i}$ denotes the set of all vertices of a simplex $T_{i} \in \mathcal{T}$ and $\operatorname{conv}(\ldots)$ denotes the convex hull of a set, then for all $T_{i}, T_{j} \in \mathcal{T}$ it is true that

$$
T_{i} \cap T_{j}=\operatorname{conv}\left(\mathcal{C}_{i} \cap \mathcal{C}_{j}\right)
$$

If a Lipschitz domain admits a triangulation, we call it a d-dimensional polyhedron. A familiy $\mathcal{F}=\left\{\mathcal{T}_{h}\right\}_{0<h \leq h_{0}}$ of triangulations is called quasi-uniform if there are positive constants $\rho_{1}$ and $\rho_{2}$ such that for all $0<h \leq h_{0}$ it holds

$$
\begin{equation*}
\max \left\{\operatorname{diam} T: T \in \mathcal{T}_{h}\right\} \leq \rho_{1} h \quad \text { and } \quad \min \left\{\operatorname{diam} B_{T}: T \in \mathcal{T}_{h}\right\} \geq \rho_{2} h \tag{2.3}
\end{equation*}
$$

Here, $B_{T}$ denotes the largest ball contained in a simplex $T$.
To approximate $(P)$, we will consider finite-dimensional minimization problems of the form

$$
\left(P_{h}\right)\left\{\begin{array}{l}
\min \frac{1}{2} a\left(v_{h}, v_{h}\right)-\left\langle f_{h}, v_{h}\right\rangle \\
\text { s.t. } v_{h} \in K_{h}:=\left\{z_{h} \in V_{h}^{0}: z_{h} \geq \psi_{h} \text { a.e. in } \Omega_{h}\right\}
\end{array} .\right.
$$

Our standing assumptions are as follows:

## Assumption 2.3.

- $\left\{\Omega_{h}\right\}_{0<h<h_{0}}$ is a family of d-dimensional polyhedra with $\Omega_{h} \subseteq \Omega$ for all $h$,
- $\left\{\mathcal{T}_{h}\right\}_{0<h \leq h_{0}}$ is a quasi-uniform family of triangulations such that each $\mathcal{T}_{h}$ is a triangulation of $\Omega_{h}$ for all $h$,
- $V_{h}^{0}:=\left\{v \in C(\bar{\Omega}):\left.v\right|_{T}\right.$ is affine for all $T \in \mathcal{T}_{h}$ and $\left.\left.v\right|_{\bar{\Omega} \backslash \Omega_{h}}=0\right\}$,
- $\psi_{h}: \Omega_{h} \rightarrow \mathbb{R}$ is measurable for all $h$,
- $f_{h} \in H^{-1}(\Omega)$ for all $h$.

Note that the condition $\Omega_{h} \subseteq \Omega$ ensures that $V_{h}^{0}$ is a subspace of $H_{0}^{1}(\Omega)$. This implies in particular that the quantities $a\left(v_{h}, v_{h}\right)$ and $\left\langle f_{h}, v_{h}\right\rangle$ are well-defined. For brevity's sake, in what follows we will often suppress the range of the mesh width $h$, i.e., we will write $\left\{\Omega_{h}\right\}$ instead of $\left\{\Omega_{h}\right\}_{0<h \leq h_{0}}, h>0$ instead of $h_{0} \geq h>0$ etc. Using again the theorem of Lions-Stampacchia, $\bar{i}$ is straightforward to prove:

THEOREM 2.4. If the admissible set $K_{h}$ is not empty, then for all $f_{h} \in H^{-1}(\Omega)$ there exists one and only one solution $u_{h} \in V_{h}^{0}$ to the problem $\left(P_{h}\right)$ and this solution is also uniquely determined by the variational inequality

$$
\begin{equation*}
u_{h} \in K_{h}: \quad a\left(u_{h}, u_{h}-v_{h}\right) \leq\left\langle f_{h}, u_{h}-v_{h}\right\rangle \quad \forall v_{h} \in K_{h} \tag{2.4}
\end{equation*}
$$

Moreover, the solution operator $S_{h}: H^{-1}(\Omega) \ni f_{h} \mapsto u_{h} \in H_{0}^{1}(\Omega)$ is Lipschitz continuous with a Lipschitz constant independent of $h$.

It should be noted that we do not assume $\psi_{h}$ to be an element of our finite element space. This will be of major importance in Section 4.
3. Discrete Supersolutions and Stability Results. To estimate the error between the continuous solution $u$ and the finite element approximation $u_{h}$, we will study the sensitivity of the solution $\operatorname{map}\left(f_{h}, \psi_{h}\right) \mapsto u_{h}$ associated with the discrete problem $\left(P_{h}\right)$. The main tool of our stability analysis will be a variant of the discrete maximum principle of Raviart-Ciarlet that is tailored to the study of the variational inequality (2.4). More precisely, we will make use of the following concept:
Definition 3.1. A function $g_{h}$ is called a discrete supersolution of the problem $\left(P_{h}\right)$ if it holds:

- $g_{h} \in V_{h}:=\left\{v \in C\left(\overline{\Omega_{h}}\right):\left.v\right|_{T}\right.$ is affine for all $\left.T \in \mathcal{T}_{h}\right\}$,
- $a\left(g_{h}, v_{h}\right) \leq\left\langle f_{h}, v_{h}\right\rangle$ for all $v_{h} \in V_{h}^{0}$ with $v_{h} \leq 0$ in $\Omega_{h}$,
- $g_{h} \geq \psi_{h}$ a.e. in $\Omega_{h}$,
- $g_{h} \geq 0$ on $\partial \Omega_{h}$.

The expression $a\left(g_{h}, v_{h}\right)$ appearing in the second point of the above definition is, of course, to be understood as

$$
a\left(g_{h}, v_{h}\right):=\int_{\Omega_{h}} \nabla g_{h} \cdot \nabla v_{h} \mathrm{~d} x
$$

In what follows, we will make frequent use of this slight abuse of notation.
Note that Definition 3.1 extends the concept of supersolutions employed in [12] straightforwardly to the discrete setting. The main idea in the following is to prove that discrete supersolutions exhibit broadly the same behavior as their continuous counterparts, i.e., to show that a discrete supersolution $g_{h}$ majorizes (at least in some
sense) the solution $u_{h}$ of the problem $\left(P_{h}\right)$. To obtain such a result, we have to restrict our analysis to triangulations of a special type:
Definition 3.2. A triangulation $\mathcal{T}_{h}$ of $\Omega_{h}$ is said to satisfy the condition $(Z)$ if

$$
\begin{equation*}
a\left(\varphi_{h}^{i}, \varphi_{h}^{j}\right)=\int_{\Omega_{h}} \nabla \varphi_{h}^{i} \cdot \nabla \varphi_{h}^{j} \mathrm{~d} x \leq 0 \quad \forall i \neq j \text { with } x_{j} \notin \partial \Omega_{h} \tag{3.1}
\end{equation*}
$$

Here, $\left\{x_{i}\right\}$ denotes the set of all vertices of the triangulation $\mathcal{T}_{h}$ (including those on the boundary $\partial \Omega_{h}$ ) and $\left\{\varphi_{h}^{i}\right\}$ denotes the nodal basis of the space $V_{h}$ (i.e., the basis with $\varphi_{h}^{i}\left(x_{l}\right)=\delta_{i l}$ for all nodes $\left.x_{l}\right)$.
The condition $(Z)$ expresses that the system matrix arising from the finite element discretization has to be a $Z$-matrix. (It is easy to see that it is even an $M$-matrix in this case). It should be noted that assumptions of the type $(Z)$ are well-known in the context of discrete maximum principles (see, e.g., [5]). In our approach, the nonnegativity condition (3.1) will come into play very naturally in the proof of Theorem 3.4. As the following lemma shows, the triangulations satisfying the condition $(Z)$ can be characterized precisely in terms of certain geometric features:

Lemma 3.3 ([19]).

- If $d=1$, then every triangulation satisfies $(Z)$.
- If $d=2$, then $(Z)$ is satisfied if and only if for each edge $E$ of $\mathcal{T}_{h}$ with $E \nsubseteq \partial \Omega_{h}$ it holds

$$
\theta_{E}^{T_{1}}+\theta_{E}^{T_{2}} \leq \pi
$$

Here, $\theta_{E}^{T_{1}}, \theta_{E}^{T_{2}} \in(0, \pi)$ denote the angles that oppose $E$ in the adjacent mesh cells $T_{1}$ and $T_{2}$ (see Figure 3.1).

- If $d>2$, then $(Z)$ is satisfied if and only if for all edges $E$ of $\mathcal{T}_{h}$ with $E \nsubseteq \partial \Omega_{h}$ it holds

$$
\sum_{T \supset E} \mathcal{H}^{d-2}\left(\kappa_{E}^{T}\right) \cot \theta_{E}^{T} \geq 0
$$

Here, for every $T=\operatorname{conv}\left(p_{1}, \ldots, p_{d+1}\right) \in \mathcal{T}_{h}$ and every $E=\operatorname{conv}\left(p_{i}, p_{j}\right) \subset T$ the quantities $\kappa_{E}^{T}$ and $\theta_{E}^{T}$ are defined by

$$
\kappa_{E}^{T}:=S_{i} \cap S_{j} \quad \text { and } \quad \theta_{E}^{T}:=\measuredangle\left(S_{i}, S_{j}\right)
$$

where $S_{i}$ and $S_{j}$ denote the $(d-1)$-dimensional simplices

$$
S_{i}:=\operatorname{conv}\left(p_{1}, \ldots, p_{i-1}, p_{i+1}, \ldots, p_{d+1}\right)
$$

and

$$
S_{j}:=\operatorname{conv}\left(p_{1}, \ldots, p_{j-1}, p_{j+1}, \ldots, p_{d+1}\right)
$$

and $\measuredangle\left(S_{i}, S_{j}\right) \in(0, \pi)$ denotes the angle enclosed by $S_{i}$ and $S_{j}$ (or the normal vectors of $S_{i}$ and $S_{j}$, to be more precise). With $\mathcal{H}^{d-2}($.$) , we mean the ( d-2$ )dimensional Hausdorff measure.


Fig. 3.1. The geometric situation in Lemma 3.3 for $d=2$ and $d=3$ (cf. [19]).
Using the results about one-sided finite element approximations found in the appendix of this paper, we can prove:

Theorem 3.4. Assume that the admissible set $K_{h}$ of the problem $\left(P_{h}\right)$ is not empty and that $\mathcal{T}_{h}$ satisfies $(Z)$. Assume further that the obstacle $\psi_{h}$ in $\left(P_{h}\right)$ satisfies

$$
\psi_{h} \in C\left(\overline{\Omega_{h}}\right) \quad \text { and }\left.\quad \psi_{h}\right|_{T} \in C^{1, \gamma}(T) \quad \forall T \in \mathcal{T}_{h}
$$

for some $\gamma \in(0,1]$ and let $\rho_{1}$ be the constant in (2.3). Then $\left(P_{h}\right)$ admits a unique solution $u_{h}$ and for every supersolution $g_{h}$ of $\left(P_{h}\right)$ it is true that

$$
\begin{equation*}
u_{h} \leq g_{h}+\frac{\sqrt{d}}{1+\gamma} \rho_{1}^{1+\gamma} h^{1+\gamma} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h}\right|_{C^{1, \gamma}(T)} \quad \text { in } \Omega_{h}, \tag{3.2}
\end{equation*}
$$

where

$$
\left|\psi_{h}\right|_{C^{1, \gamma}(T)}:=\max _{k=1, \ldots, d} \sup _{x \neq y \in T} \frac{\left|\partial_{k} \psi_{h}(x)-\partial_{k} \psi_{h}(y)\right|}{\|x-y\|^{\gamma}} .
$$

The above theorem shows that discrete supersolutions at least approximately behave as expected: They are larger than the solution $u_{h}$ modulo an error that depends on the mesh width $h$ and the curvature of the obstacle $\psi_{h}$. The inequality $u_{h} \leq g_{h}$, i.e., the behavior observed in the continuous setting (cf. [12, Theorem II6.4]), is only obtained if $\psi_{h}$ is an element of the space $V_{h}$.
Proof of Theorem 3.4. The existence of the solution $u_{h}$ follows straightforwardly from Theorem 2.4. To prove inequality (3.2), we will use an argument similar to that employed in the continuous setting (cf. [12]). In a first step, we define $g_{h}^{\prime}:=g_{h}+C$, where

$$
C:=\frac{\sqrt{d}}{1+\gamma} \rho_{1}^{1+\gamma} h^{1+\gamma} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h}\right|_{C^{1, \gamma}(T)} .
$$

Note that $g_{h}^{\prime}$ is again a supersolution since the addition of a positive constant to the function $g_{h}$ has no effect on the properties in Definition 3.1. We now consider the unique element $v_{h}$ in $V_{h}$ with

$$
\begin{equation*}
v_{h}\left(x_{i}\right)=\min \left(u_{h}\left(x_{i}\right), g_{h}^{\prime}\left(x_{i}\right)\right) \quad \forall x_{i}, \tag{3.3}
\end{equation*}
$$

where $\left\{x_{i}\right\}$ again denotes the set off all vertices of $\mathcal{T}_{h}$ (including those on the boundary $\partial \Omega_{h}$ ). Since $u_{h} \in V_{h}^{0}$, the function $v_{h}$ can be identified with an element of $V_{h}^{0}$ and
from (3.3) we readily obtain that $v_{h} \leq u_{h}$ holds everywhere in $\Omega_{h}$. Furthermore, it follows from our construction that $v_{h} \geq \psi_{h}$. To see this, we use that according to Theorem A. 4 from the appendix (applied to $z:=u_{h}-\psi_{h}$ ), for every mesh cell $T \in \mathcal{T}_{h}$ with vertices $p_{1}, \ldots, p_{d+1}$ we can find an affine linear function $\psi_{h}^{T}$ on $T$ such that $\psi_{h} \leq \psi_{h}^{T} \leq u_{h}$ holds on $T$ and such that

$$
0 \leq \psi_{h}^{T}\left(p_{k}\right)-\psi_{h}\left(p_{k}\right) \leq \frac{\sqrt{d}}{1+\gamma} \rho_{1}^{1+\gamma} h^{1+\gamma}\left|\psi_{h}\right|_{C^{1, \gamma}(T)} \quad \forall k=1, . ., d+1
$$

This yields that $v_{h}$ satisfies

$$
v_{h}\left(p_{k}\right)=\min \left(u_{h}\left(p_{k}\right), g_{h}\left(p_{k}\right)+C\right) \geq \min \left(\psi_{h}^{T}\left(p_{k}\right), \psi_{h}\left(p_{k}\right)+C\right) \geq \psi_{h}^{T}\left(p_{k}\right)
$$

for all $k=1, \ldots, d+1$. From the affine linearity of $v_{h}$ and $\psi_{h}^{T}$ on $T$, it now follows $v_{h} \geq \psi_{h}^{T} \geq \psi_{h}$ which implies $v_{h} \geq \psi_{h}$ on $\Omega_{h}$ as claimed. From the second property in Definition 3.1 and the variational inequality (2.4), we may now deduce:

$$
a\left(g_{h}^{\prime}, v_{h}-u_{h}\right) \leq\left\langle f_{h}, v_{h}-u_{h}\right\rangle \quad \text { and } \quad a\left(u_{h}, u_{h}-v_{h}\right) \leq\left\langle f_{h}, u_{h}-v_{h}\right\rangle .
$$

If we add these inequalities and define $y_{i}:=u_{h}\left(x_{i}\right)-g_{h}^{\prime}\left(x_{i}\right)$ for all nodes $x_{i}$, we obtain (using the properties of $v_{h}$ )

$$
\begin{align*}
0 & \geq a\left(u_{h}-g_{h}^{\prime}, u_{h}-v_{h}\right) \\
& =\sum_{x_{i}} y_{i} \max \left(0, y_{i}\right) a\left(\varphi_{h}^{i}, \varphi_{h}^{i}\right)+\sum_{x_{i} \neq x_{j}} y_{i} \max \left(0, y_{j}\right) a\left(\varphi_{h}^{i}, \varphi_{h}^{j}\right) \\
& =\sum_{x_{i}} \max \left(0, y_{i}\right)^{2} a\left(\varphi_{h}^{i}, \varphi_{h}^{i}\right)+\sum_{x_{i} \neq x_{j} \text { and } x_{j} \notin \partial \Omega_{h}} y_{i} \max \left(0, y_{j}\right) a\left(\varphi_{h}^{i}, \varphi_{h}^{j}\right), \tag{3.4}
\end{align*}
$$

where $\left\{\varphi_{h}^{i}\right\}$ again denotes the nodal basis of $V_{h}$. Because of the condition $(Z)$, however, we also know that for all $i, j$ with $x_{i} \neq x_{j}$ and $x_{j} \notin \partial \Omega_{h}$ it holds

$$
y_{i} \max \left(0, y_{j}\right) a\left(\varphi_{h}^{i}, \varphi_{h}^{j}\right) \geq \max \left(0, y_{i}\right) \max \left(0, y_{j}\right) a\left(\varphi_{h}^{i}, \varphi_{h}^{j}\right)
$$

Thus, (3.4) implies

$$
0 \geq \sum_{x_{i}} \sum_{x_{j}} \max \left(0, y_{i}\right) \max \left(0, y_{j}\right) a\left(\varphi_{h}^{i}, \varphi_{h}^{j}\right)=a\left(u_{h}-v_{h}, u_{h}-v_{h}\right) \geq 0
$$

and consequently

$$
u_{h}\left(x_{i}\right)-v_{h}\left(x_{i}\right)=\max \left(0, u_{h}\left(x_{i}\right)-g_{h}^{\prime}\left(x_{i}\right)\right)=0 \quad \forall x_{i}
$$

Using again the piecewise linearity of the involved functions, we may deduce

$$
u_{h} \leq g_{h}^{\prime}=g_{h}+C=g_{h}+\frac{\sqrt{d}}{1+\gamma} \rho_{1}^{1+\gamma} h^{1+\gamma} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h}\right|_{C^{1, \gamma}(T)} \quad \text { in } \Omega_{h} .
$$

This completes the proof.
Theorem 3.4 allows to analyze the sensitivity of the solution $u_{h}$ with respect to perturbations of the obstacle $\psi_{h}$ and the force $f_{h}$ :

Theorem 3.5. Consider two discrete obstacle problems of the form

$$
\left(P_{h, i}\right)\left\{\begin{array}{l}
\min \frac{1}{2} a\left(v_{h}, v_{h}\right)-\left\langle f_{h, i}, v_{h}\right\rangle \\
\text { s.t. } v_{h} \in K_{h, i}:=\left\{z_{h} \in V_{h}^{0}: z_{h} \geq \psi_{h, i} \text { a.e. in } \Omega_{h}\right\}
\end{array}, \quad i=1,2,\right.
$$

and assume that:

- $f_{h, 1}, f_{h, 2} \in H^{-1}(\Omega)$,
- the underlying triangulation $\mathcal{T}_{h}$ satisfies $(Z)$,
- $\psi_{h, 1}, \psi_{h, 2} \in C\left(\overline{\Omega_{h}}\right)$ and $K_{h, 1} \neq \emptyset, K_{h, 2} \neq \emptyset$,
- there exist $\gamma_{1}, \gamma_{2} \in(0,1]$ such that $\left.\psi_{h, i}\right|_{T} \in C^{1, \gamma_{i}}(T)$ for all $T \in \mathcal{T}_{h}, i=1,2$.

Let $\rho_{1}$ be the constant in (2.3). Then $\left(P_{h, 1}\right)$ and $\left(P_{h, 2}\right)$ admit unique solutions $u_{h, 1}$ and $u_{h, 2}$ and there exists a constant $C>0$ independent of $h$ such that

$$
\begin{align*}
&\left\|\left(u_{h, 1}-u_{h, 2}\right)^{+}\right\|_{L^{\infty}} \\
& \leq\left\|\left(\psi_{h, 1}-\psi_{h, 2}\right)^{+}\right\|_{L^{\infty}}+C r(h)\left\|f_{h, 1}-f_{h, 2}\right\|_{H^{-1}} \\
&+\frac{\sqrt{d}}{1+\gamma_{1}}\left(\rho_{1} h\right)^{1+\gamma_{1}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h, 1}\right|_{C^{1, \gamma_{1}}(T)} \tag{3.5}
\end{align*}
$$

and

$$
\begin{align*}
&\left\|\left(u_{h, 1}-u_{h, 2}\right)^{-}\right\|_{L^{\infty}} \\
& \leq\left\|\left(\psi_{h, 1}-\psi_{h, 2}\right)^{-}\right\|_{L^{\infty}}+C r(h)\left\|f_{h, 1}-f_{h, 2}\right\|_{H^{-1}} \\
&+\frac{\sqrt{d}}{1+\gamma_{2}}\left(\rho_{1} h\right)^{1+\gamma_{2}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h, 2}\right|_{C^{1, \gamma_{2}}(T)} \tag{3.6}
\end{align*}
$$

Here, $v^{+}:=\max (0, v)$ and $v^{-}:=\min (0, v)$ denote the positive and the negative part of a function, respectively, and $r(h)$ is defined by:

$$
r(h):= \begin{cases}1 & \text { if } d=1  \tag{3.7}\\ (1+|\log h|)^{1 / 2} & \text { if } d=2 . \\ h^{1-d / 2} & \text { if } d \geq 3\end{cases}
$$

Proof. The unique solvability of the problems $\left(P_{h, 1}\right)$ and $\left(P_{h, 2}\right)$ is a straightforward consequence of Theorem 2.4. It remains to prove the estimates (3.5) and (3.6). If we assume first that $f_{h, 1}=f_{h, 2}$ and define $g_{h, 1}:=u_{h, 2}+\left\|\left(\psi_{h, 1}-\psi_{h, 2}\right)^{+}\right\|_{L^{\infty}}$, then it certainly holds $g_{h, 1} \in V_{h}, g_{h, 1} \geq 0$ on $\partial \Omega_{h}$ and

$$
g_{h, 1} \geq u_{h, 2}+\psi_{h, 1}-\psi_{h, 2} \geq \psi_{h, 1} \text { in } \Omega_{h}
$$

From the variational inequality associated with $\left(P_{h, 2}\right)$ and the definition of $a(.,$.$) , it$ follows further that $g_{h, 1}$ satisfies

$$
a\left(g_{h, 1}, v_{h}\right)=a\left(u_{h, 2}, v_{h}\right)=a\left(u_{h, 2}, u_{h, 2}-\left(u_{h, 2}-v_{h}\right)\right) \leq\left\langle f_{h, 2}, v_{h}\right\rangle=\left\langle f_{h, 1}, v_{h}\right\rangle
$$

for every $v_{h} \in V_{h}^{0}$ with $v_{h} \leq 0$ in $\Omega_{h}$. Thus, $g_{h, 1}$ is a supersolution for $\left(P_{h, 1}\right)$ and we may deduce from our last theorem that

$$
u_{h, 1} \leq g_{h, 1}+\frac{\sqrt{d}}{1+\gamma_{1}}\left(\rho_{1} h\right)^{1+\gamma_{1}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h, 1}\right|_{C^{1, \gamma_{1}}(T)} \quad \text { in } \Omega_{h}
$$

which implies

$$
\left\|\left(u_{h, 1}-u_{h, 2}\right)^{+}\right\|_{L^{\infty}} \leq\left\|\left(\psi_{h, 1}-\psi_{h, 2}\right)^{+}\right\|_{L^{\infty}}+\frac{\sqrt{d}}{1+\gamma_{1}}\left(\rho_{1} h\right)^{1+\gamma_{1}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h, 1}\right|_{C^{1, \gamma_{1}}(T)} .
$$

This is exactly (3.5) for $f_{h, 1}=f_{h, 2}$. The inequality (3.6) is obtained analogously if we interchange the roles of $u_{h, 1}$ and $u_{h, 2}$. This proves the claim for discrete obstacle problems with identical forces. Assume now that $f_{h, 1} \neq f_{h, 2}$ and denote with $u_{h, i, j}$ the solution of the discrete obstacle problem with obstacle $\psi_{h, i}$ and force $f_{h, j}$, then it follows from the triangle inequality, the Lipschitz property in Theorem 2.4 and well-known inverse estimates (see, e.g., [3, Section 4.5, 4.9]) that

$$
\begin{aligned}
& \left\|\left(u_{h, 1}-u_{h, 2}\right)^{+}\right\|_{L^{\infty}} \\
& \quad \leq\left\|\left(u_{h, 1,1}-u_{h, 2,1}\right)^{+}\right\|_{L^{\infty}}+\left\|u_{h, 2,1}-u_{h, 2,2}\right\|_{L^{\infty}} \\
& \leq\left\|\left(u_{h, 1,1}-u_{h, 2,1}\right)^{+}\right\|_{L^{\infty}}+C_{1} r(h)\left\|u_{h, 2,1}-u_{h, 2,2}\right\|_{H^{1}} \\
& \leq\left\|\left(\psi_{h, 1}-\psi_{h, 2}\right)^{+}\right\|_{L^{\infty}}+C_{2} r(h)\left\|f_{h, 1}-f_{h, 2}\right\|_{H^{-1}} \\
& \quad+\frac{\sqrt{d}}{1+\gamma_{1}} \rho_{1}^{1+\gamma_{1}} h^{1+\gamma_{1}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h, 1}\right|_{C^{1, \gamma_{1}}(T)}
\end{aligned}
$$

where $C_{1}, C_{2}$ are constants independent of $h$. This proves (3.5) in the general case. The estimate (3.6) is again obtained by interchanging roles.
As the above result shows, for fixed $f_{h}$ the solution operator $\psi_{h} \mapsto u_{h}$ of the problem $\left(P_{h}\right)$ is not Lipschitz continuous as a function from (a subset of) $L^{\infty}(\Omega)$ to $L^{\infty}(\Omega)$. We only obtain a Lipschitz-like estimate with an error that again depends on the mesh width $h$ and the curvature of the involved obstacles. This is a major difference to the continuous setting where it can be shown easily that the solutions $u_{1}$ and $u_{2}$ of two obstacle problems with $L^{\infty}$-obstacles $\psi_{1}$ and $\psi_{2}$ and identical forces satisfy $\left\|u_{1}-u_{2}\right\|_{L^{\infty}} \leq C\left\|\psi_{1}-\psi_{2}\right\|_{L^{\infty}}$ (cf. [12, Theorem IV8.5]). It should be noted that neither the continuous solution $u$ nor the obstacle $\psi$ or the domain $\Omega$ have been relevant for the derivation of (3.5) and (3.6). Up to now, we have solely worked with the discrete problems.
4. $\mathbf{L}^{\infty}$-Error Estimates. A priori estimates for the error $\left\|u-u_{h}\right\|_{L^{\infty}}$ can be derived straightforwardly from Theorem 3.5. We just have to observe the following:
Lemma 4.1. If $u \in H_{0}^{1}(\Omega)$ is the solution to the obstacle problem $(P)$ and $R_{h} u$ the Ritz projection of $u$, i.e., the unique element of $V_{h}^{0}$ satisfying

$$
\begin{equation*}
a\left(R_{h} u, v_{h}\right)=a\left(u, v_{h}\right) \quad \forall v_{h} \in V_{h}^{0} \tag{4.1}
\end{equation*}
$$

then $R_{h} u$ is also the unique solution to the discrete obstacle problem

$$
\left(Q_{h}\right)\left\{\begin{array}{l}
\min \frac{1}{2} a\left(v_{h}, v_{h}\right)-\left\langle f, v_{h}\right\rangle \\
\text { s.t. } v_{h} \in V_{h}^{0} \text { and } v_{h} \geq \psi+R_{h} u-u \text { a.e. in } \Omega_{h}
\end{array} .\right.
$$

Recall that we have assumed $\Omega_{h} \subseteq \Omega$ (see Assumption 2.3). This ensures that $u$ and $\psi$ are defined everywhere in $\Omega_{h}$ and that the constraint in $\left(Q_{h}\right)$ makes sense.
Proof of Lemma 4.1. The Ritz projection $R_{h} u$ is obviously admissible for $\left(Q_{h}\right)$ and because of (4.1) and the variational inequality (2.1), it holds

$$
a\left(R_{h} u, R_{h} u-v_{h}\right)=a\left(u, R_{h} u-v_{h}\right)=a\left(u, u-\left(u-R_{h} u+v_{h}\right)\right) \leq\left\langle f, R_{h} u-v_{h}\right\rangle
$$

for all $v_{h} \in V_{h}^{0}$ with $v_{h} \geq \psi+R_{h} u-u$, i.e., $u-R_{h} u+v_{h} \geq \psi$. This shows that $R_{h} u$ is indeed the solution to $\left(Q_{h}\right)$ and completes the proof (cf. Theorem 2.4).

Note that Lemma 4.1 holds without any further assumptions on the regularity of the functions $u$ and $\psi$ and that the obstacle $\psi+R_{h} u-u$ appearing in $\left(Q_{h}\right)$ is typically not piecewise linear. By applying Theorem 3.5 to $\left(Q_{h}\right)$ and the problem $\left(P_{h}\right)$ used for the finite element approximation, we obtain:

THEOREM 4.2. Assume that $(P)$ admits a solution $u$ and denote with $R_{h} u$ the Ritz projection of $u$ as defined in (4.1). Suppose further that the following is satisfied:

- $u, \psi \in C(\bar{\Omega}), \psi_{h} \in C\left(\overline{\Omega_{h}}\right)$ and $K_{h} \neq \emptyset$,
- $\exists \gamma_{1}, \gamma_{2} \in(0,1]$ with $\left.\psi_{h}\right|_{T} \in C^{1, \gamma_{1}}(T)$ and $\left.u\right|_{T},\left.\psi\right|_{T} \in C^{1, \gamma_{2}}(T)$ for all $T \in \mathcal{T}_{h}$,
- the triangulation $\mathcal{T}_{h}$ satisfies $(Z)$.

Then $\left(P_{h}\right)$ admits a unique solution $u_{h}$ and there exists a constant $C>0$ independent of $h$ such that

$$
\begin{align*}
\|(u- & \left.u_{h}\right)^{-} \|_{L^{\infty}\left(\Omega_{h}\right)} \\
\leq & \left\|\left(u-R_{h} u\right)^{-}\right\|_{L^{\infty}\left(\Omega_{h}\right)}+\left\|\left(\psi_{h}-\psi+u-R_{h} u\right)^{+}\right\|_{L^{\infty}\left(\Omega_{h}\right)} \\
& +C r(h)\left\|f-f_{h}\right\|_{H^{-1}(\Omega)}+\frac{\sqrt{d}}{1+\gamma_{1}}\left(\rho_{1} h\right)^{1+\gamma_{1}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h}\right|_{C^{1, \gamma_{1}}(T)} \tag{4.2}
\end{align*}
$$

and

$$
\begin{align*}
\|(u- & \left.u_{h}\right)^{+} \|_{L^{\infty}\left(\Omega_{h}\right)} \\
\leq & \left\|\left(u-R_{h} u\right)^{+}\right\|_{L^{\infty}\left(\Omega_{h}\right)}+\left\|\left(\psi_{h}-\psi+u-R_{h} u\right)^{-}\right\|_{L^{\infty}\left(\Omega_{h}\right)} \\
& +C r(h)\left\|f-f_{h}\right\|_{H^{-1}(\Omega)}+\frac{\sqrt{d}}{1+\gamma_{2}}\left(\rho_{1} h\right)^{1+\gamma_{2}} \max _{T \in \mathcal{T}_{h}}|\psi-u|_{C^{1, \gamma_{2}}(T)} \tag{4.3}
\end{align*}
$$

Here, $r(h)$ is again defined by (3.7).
With the above theorem we have reduced the problem of finding an a priori estimate for the error $\left\|u-u_{h}\right\|_{L^{\infty}}$ to that of estimating the $L^{\infty}$-error between $u$ and the Ritz projection $R_{h} u$. The approximation properties of $R_{h} u$, however, have been studied by numerous authors and estimates for the quantity $\left\|u-R_{h} u\right\|_{L^{\infty}\left(\Omega_{h}\right)}$ are well-known. The following result can be found, for example, in [17]:

Lemma 4.3. Assume that $\partial \Omega$ is smooth, that $u \in H_{0}^{1}(\Omega) \cap C(\bar{\Omega})$ and that there exists a constant $\delta>0$ independent of $h$ such that

$$
\max _{x \in \partial \Omega_{h}} \operatorname{dist}(x, \partial \Omega) \leq \delta h^{2}
$$

Then there exists a constant $C>0$ independent of $h$ such that

$$
\left\|u-R_{h} u\right\|_{L^{\infty}\left(\Omega_{h}\right)} \leq C|\log h|^{\alpha} \inf _{v_{h} \in V_{h}^{0}}\left\|u-v_{h}\right\|_{L^{\infty}\left(\Omega_{h}\right)}
$$

with $\alpha=0$ for $d=1$ and $\alpha=1$ for $d>1$.
Combining Theorem 4.2 and Lemma 4.3 yields:

Corollary 4.4. Let the assumptions of Theorem 4.2 and Lemma 4.3 hold. Then there exists a constant $C>0$ independent of $h$ such that

$$
\begin{align*}
\|(u- & \left.u_{h}\right)^{-} \|_{L^{\infty}\left(\Omega_{h}\right)} \\
\leq & \left\|\left(\psi-\psi_{h}\right)^{-}\right\|_{L^{\infty}\left(\Omega_{h}\right)}+\frac{\sqrt{d}}{1+\gamma_{1}}\left(\rho_{1} h\right)^{1+\gamma_{1}} \max _{T \in \mathcal{T}_{h}}\left|\psi_{h}\right|_{C^{1, \gamma_{1}}(T)} \\
& +C r(h)\left\|f-f_{h}\right\|_{H^{-1}(\Omega)}+C|\log h|^{\alpha} \inf _{v_{h} \in V_{h}^{0}}\left\|u-v_{h}\right\|_{L^{\infty}\left(\Omega_{h}\right)} \tag{4.4}
\end{align*}
$$

and

$$
\begin{align*}
& \left\|\left(u-u_{h}\right)^{+}\right\|_{L^{\infty}\left(\Omega_{h}\right)} \\
& \leq \leq\left\|\left(\psi-\psi_{h}\right)^{+}\right\|_{L^{\infty}\left(\Omega_{h}\right)}+\frac{\sqrt{d}}{1+\gamma_{2}}\left(\rho_{1} h\right)^{1+\gamma_{2}} \max _{T \in \mathcal{T}_{h}}|\psi-u|_{C^{1, \gamma_{2}}(T)} \\
& \quad+C r(h)\left\|f-f_{h}\right\|_{H^{-1}(\Omega)}+C|\log h|^{\alpha} \inf _{v_{h} \in V_{h}^{0}}\left\|u-v_{h}\right\|_{L^{\infty}\left(\Omega_{h}\right)}, \tag{4.5}
\end{align*}
$$

where $\alpha$ and $r(h)$ are defined as before.
As a consequence of Corollary 4.4, we obtain in particular:
Corollary 4.5. Assume that:

- $\partial \Omega$ is smooth,
- $f \in L^{q}(\Omega)$ and $\psi \in W^{2, q}(\Omega)$ for some $\max (d, 2)<q<\infty$,
$-\operatorname{tr} \psi \leq 0$ a.e. on $\partial \Omega$,
- there exists a constant $\delta>0$ independent of $h$ such that

$$
\max _{x \in \partial \Omega_{h}} \operatorname{dist}(x, \partial \Omega) \leq \delta h^{2}
$$

- the triangulation $\mathcal{T}_{h}$ satisfies $(Z)$.

Suppose further that one of the following holds:
a) $\psi_{h}$ is equal to the Lagrange interpolant $I_{h} \psi \in V_{h}$ of $\psi$ and $K_{h} \neq \emptyset$.
b) $\psi_{h}$ is equal to the restriction $\left.\psi\right|_{\Omega_{h}}$ and $K_{h} \neq \emptyset$.

Then $(P)$ and $\left(P_{h}\right)$ admit unique solutions $u$ and $u_{h}$, it holds $u \in H_{0}^{1}(\Omega) \cap W^{2, q}(\Omega)$ and there exists a constant $C>0$ independent of $h$ such that

$$
\begin{align*}
& \left\|u-u_{h}\right\|_{L^{\infty}\left(\Omega_{h}\right)} \\
& \quad \leq C|\log h|^{\alpha} h^{2-d / q}\left(\|f\|_{L^{q}(\Omega)}+\|\psi\|_{W^{2, q}(\Omega)}\right)+C r(h)\left\|f-f_{h}\right\|_{H^{-1}(\Omega)}, \tag{4.6}
\end{align*}
$$

where $\alpha$ and $r(h)$ are defined as before.
Proof. The unique solvability of the problems $(P)$ and $\left(P_{h}\right)$ and the $W^{2, q}$-regularity of the solution $u$ are direct consequences of Theorem 2.1 and Theorem 2.4. The error estimate (4.6) follows straightforwardly from (4.4) and (4.5). We just have to employ standard results about the accuracy of the Lagrange interpolant (as found, e.g., in [3, Theorem 4.4.20]) and the embedding $W^{2, q}(\Omega) \hookrightarrow C^{1,1-d / q}(\bar{\Omega})$.
Note that in case $a$ ), (4.6) is the 'standard' $L^{\infty}$-error estimate for the obstacle problem that is usually found in the literature (cf. [2,13,15]).
5. Concluding Remarks. The method that we have employed in the last two sections to derive a priori error estimates for the obstacle problem $(P)$ has some advantages that we would like to point out here:
First of all, our approach is more flexible than the traditional one since we do not require $\psi_{h}$ to be the Lagrange interpolant $I_{h} \psi$ of the continuous obstacle $\psi$ (or an element of the finite element space at all, cf. Theorem 4.2). Moreover, we can treat the case $\Omega_{h} \subset \Omega$ with ease since the relation between the domains $\Omega$ and $\Omega_{h}$ is completely irrelevant for the stability analysis that our proofs are based on (cf. Theorem 3.5).
Second, our results provide slightly more information about the behavior of the approximation error than those found in the literature. We obtain, for example, in a natural way separate estimates for the quantities $\left(u-u_{h}\right)^{+}$and $\left(u-u_{h}\right)^{-}$that allow to study in greater detail how the accuracy of the finite element method is affected by the choice of $\psi_{h}$ (cf. (4.2) and (4.3)).

Lastly, our approach demonstrates that the problem of estimating the error between the continuous solution $u$ and the finite element approximation $u_{h}$ can be identified with a problem of sensitivity analysis: If we know how the solution of the discrete obstacle problem $\left(Q_{h}\right)$ changes when the obstacle $\psi+R_{h} u-u$ is replaced with $\psi_{h}$, then we also know how the quantities $u-u_{h}$ and $u-R_{h} u$, i.e., the errors associated with the constraint and the unconstraint setting, are related to each other and vice versa. Note that this interpretation is only possible when the obstacles in the discrete problems are allowed to be arbitrary measurable functions (cf. the definition of $\left(Q_{h}\right)$ ).
The above perspective on the a priori error analysis turns out to be very advantageous when error estimates in lower $L^{p}$-norms are considered. In [4], for example, it was used to construct two counterexamples which demonstrate (among other things) that the estimate (4.6) is optimal in the one-dimensional case in the sense that there exist situations where the assumptions of Corollary 4.5 are satisfied and where it holds $\left\|u-u_{h}\right\|_{L^{p}\left(\Omega_{h}\right)}=\operatorname{ord}\left(h^{2-1 / q}\right)$ for all $1 \leq p \leq \infty$. Interestingly, the latter is true regardless of whether the Lagrange interpolant $I_{h} \psi$ or the restriction $\left.\psi\right|_{\Omega_{h}}$ is chosen as $\psi_{h}$ in $\left(P_{h}\right)$. We refer to [4] for a detailed discussion of this topic.
It should be noted that the situation is much less clear in higher dimensions and that it is (at least to the author's best knowledge) presently unknown if an $L^{p}$-error estimate of the form $\left\|u-u_{h}\right\|_{L^{p}\left(\Omega_{h}\right)}=\mathcal{O}\left(h^{\gamma}\right)$ with $\gamma>2-d / q$ and $\gamma>1$ can be obtained for an obstacle problem with $u, \psi \in W^{2, q}(\Omega)$ if the dimension is greater than one. A further open question is whether the condition $(Z)$ can be weakened. The results found in [6] indicate that the latter might be the case and that it might be sufficient to assume that (3.1) holds in an appropriately chosen subset of $\Omega_{h}$ to derive Theorem 3.4 (cf. also the results in [16]). A proof of this conjecture, however, is still pending.

## Appendix A. Finite Element Approximation under Inequality Constraints.

In this section, we prove the approximation results that we have used in the proof of Theorem 3.4. We will be mainly concerned with the following task:

$$
\begin{aligned}
& \text { Assume that } T \text { is a closed d-dimensional simplex with vertices } p_{1}, \ldots, p_{d+1} \\
& \text { and let } z \text { be a function satisfying } 0 \leq z \in C^{1, \gamma}(T) \text { for some } 0<\gamma \leq 1 \text {. } \\
& \text { Find an affine linear function } z_{T} \text { such that } 0 \leq z_{T} \leq z \text { holds in } T \\
& \text { and such that } z_{T} \text { approximates } z \text { as accurately as possible. }
\end{aligned}
$$

The above approximation problem has already been studied by Mosco in [14] and Strang in [18] for one- and two-dimensional $H^{2}$-functions. The method of proof that we will employ in this section is closely related to the approach of these two authors. Our analysis, however, also covers the higher-dimensional case.

To construct an approximation $z_{T}$ with the desired properties, we introduce a partial order on the set of affine linear functions on $T$ (analogously to [14]) and define:
Definition A.1. Let $T \subset \mathbb{R}^{d}$ and $0 \leq z \in C^{1, \gamma}(T), 0<\gamma \leq 1$, be as above and let

$$
V:=\{v: T \rightarrow \mathbb{R}: v \text { affine linear with } 0 \leq v \leq z \text { in } T\}
$$

Then a function $v \in V$ is called a maximal element of $V$ if for every affine linear function $w$ with $w \not \equiv 0$ and $w \geq 0$ in $T$ it holds $v+w \notin V$.
Using standard arguments, it is easy to prove:
Lemma A.2. The set $V$ always admits at least one maximal element.
Proof. If we denote with $p_{1}, \ldots, p_{d+1}$ the vertices of $T$ and define

$$
U=\left\{v \in \mathbb{R}^{d+1}: 0 \leq \sum_{i=1}^{d+1} \lambda_{i} v_{i} \leq z\left(\sum_{i=1}^{d+1} \lambda_{i} p_{i}\right) \text { for all } \lambda_{i} \geq 0 \text { with } \sum_{i=1}^{d+1} \lambda_{i}=1\right\}
$$

then $U$ is closed, non-empty and bounded. Thus, the function $f(v):=\sum_{i=1}^{d+1} v_{i}$ attains its supremum in $U$ in some $v$. Because of its maximality, this $v$ has to satisfy $v+w \notin U$ for all $w \in \mathbb{R}^{d+1} \backslash\{0\}$ with $w \geq 0$ (componentwise). This, however, implies that the affine linear map

$$
x=\sum_{i=1}^{d+1} \lambda_{i} p_{i} \in T \mapsto \sum_{i=1}^{d+1} \lambda_{i} v_{i}
$$

defined in the barycentric coordinates w.r.t. the vertices $p_{i}$ of the simplex $T$ is a maximal element of $V$.

To estimate the difference between $z$ and a maximal element of $V$, we observe the following:
Lemma A.3. Let $T \subset \mathbb{R}^{d}$ and $0 \leq z \in C^{1, \gamma}(T), 0<\gamma \leq 1$, be as above and denote with $p_{1}, \ldots, p_{d+1}$ the vertices of $T$. Let $v$ be a maximal element of $V$ and define

$$
E(v):=\{\zeta \in T: z(\zeta)=v(\zeta)\}
$$

Then $E(v)$ is not empty and if it holds $E(v) \subseteq \operatorname{conv}\left(p_{1}, \ldots, p_{k-1}, p_{k+1}, \ldots, p_{d+1}\right)$, where $\operatorname{conv}(\ldots)$ denotes the convex hull, then there exists a $\zeta \in E(v)$ such that

$$
\begin{equation*}
\nabla(z-v)(\zeta) \cdot\left(p_{k}-\zeta\right)=0 \tag{A.1}
\end{equation*}
$$

Proof. The non-emptiness of $E(v)$ follows trivially from the maximality of $v$. To prove the second part of the lemma, we assume w.l.o.g. that $k=d+1$, that $\operatorname{conv}\left(p_{1}, \ldots, p_{d}\right) \subset$ $\mathbb{R}^{d-1} \times\{0\}$, and that $p_{d+1} \in \mathbb{R}^{d-1} \times(0, \infty)$. Since all $\zeta \in E(v)$ are global minima of the function $z-v \in C^{1}(T)$, it necessarily holds

$$
\begin{equation*}
D(\zeta):=\nabla(z-v)(\zeta) \cdot\left(p_{d+1}-\zeta\right) \geq 0 \quad \forall \zeta \in E(v) \tag{A.2}
\end{equation*}
$$

From the compactness of $E(v)$ and the continuity of the function $D: T \rightarrow \mathbb{R}$, we obtain that there exists a $\zeta^{\prime} \in E(v)$ with

$$
\begin{equation*}
D\left(\zeta^{\prime}\right)=\min _{\zeta \in E(v)} D(\zeta)=: m \geq 0 \tag{A.3}
\end{equation*}
$$

If the minimum $m$ is zero, then the claim is obviously true. If $m>0$, then there exists an open ball $B(\zeta)$ around each $\zeta \in E(v)$ such that $D(x)>m / 2$ holds for all $x \in T \cap B(\zeta)$ and we may define

$$
\begin{aligned}
& B^{\prime}(\zeta) \\
& \quad:=B(\zeta) \cap\left\{\lambda x+(1-\lambda) p_{d+1}: x \in B(\zeta) \cap T \cap\left(\mathbb{R}^{d-1} \times\{0\}\right) \text { and } \lambda \in(0,1]\right\}
\end{aligned}
$$

and

$$
E^{\prime}(v):=\bigcup_{\zeta \in E(v)} B^{\prime}(\zeta)
$$

Note that it follows from our construction that $E(v) \subset E^{\prime}(v) \subset T$ (cf. Figure A.1). Moreover, $E^{\prime}(v)$ is relatively open in $T$ and it holds

$$
\nabla(z-v)(x) \cdot\left(p_{d+1}-x\right)>\frac{1}{2} m \quad \forall x \in E^{\prime}(v)
$$

Suppose now that $c$ is a constant satisfying $0<c<c^{\prime}:=m /\left(4 x_{d}\left(p_{d+1}\right)\right)$, where $x_{d}\left(p_{d+1}\right)>0$ denotes the $d$-th coordinate of the point $p_{d+1}$, and consider the function $v_{c}(x):=v(x)+c x_{d}$. Then $v_{c}$ is obviously affine and it holds

$$
\nabla\left(z-v_{c}\right)(x) \cdot\left(p_{d+1}-x\right) \geq \frac{1}{2} m-c x_{d}\left(p_{d+1}\right) \geq \frac{1}{4} m \quad \forall x \in E^{\prime}(v)
$$

From $z-v_{c}=z-v \geq 0$ on $T \cap\left(\mathbb{R}^{d-1} \times\{0\}\right)$, the mean value theorem, and the fact that for every $x \in E^{\prime}(v)$ the line between $x$ and the unique $x^{\prime}=x^{\prime}(x)$ with $x^{\prime} \in T \cap\left(\mathbb{R}^{d-1} \times\{0\}\right)$ and $x \in \operatorname{conv}\left(x^{\prime}, p_{d+1}\right)$ is contained in $E^{\prime}(v)$ (cf. Figure A.1), it follows further

$$
\begin{aligned}
\left(z-v_{c}\right)(x) & \geq\left(z-v_{c}\right)(x)-\left(z-v_{c}\right)\left(x^{\prime}\right) \\
& =\int_{0}^{1} \nabla\left(z-v_{c}\right)\left(x^{\prime}+t\left(x-x^{\prime}\right)\right) \cdot\left(x-x^{\prime}\right) \mathrm{d} t \\
& =\int_{0}^{1} \nabla\left(z-v_{c}\right)\left(x^{\prime}+t\left(x-x^{\prime}\right)\right) \cdot \frac{\left(p_{d+1}-\left(x^{\prime}+t\left(x-x^{\prime}\right)\right)\right.}{\| p_{d+1}-\left(x^{\prime}+t\left(x-x^{\prime}\right) \|\right.}\left\|x-x^{\prime}\right\| \mathrm{d} t \\
& \geq \frac{1}{4} m \int_{0}^{1} \frac{\left\|x-x^{\prime}\right\|}{\| p_{d+1}-\left(x^{\prime}+t\left(x-x^{\prime}\right) \|\right.} \mathrm{d} t \\
& \geq 0 \quad \forall x \in E^{\prime}(v) .
\end{aligned}
$$

To avoid a contradiction with the maximality of $v$, it now has to hold that for every $0<c<c^{\prime}$ there exists at least one $x_{c} \in T \backslash E^{\prime}(v)$ with $\left(z-v_{c}\right)\left(x_{c}\right)<0$. The set $T \backslash E^{\prime}(v)$, however, is compact. This implies that we can find a sequence $c_{n} \rightarrow 0$ such that $x_{c_{n}}$ converges to an $x_{0} \in T \backslash E^{\prime}(v)$ and such a limit $x_{0}$ has to satisfy

$$
0 \geq \lim _{c_{n} \rightarrow 0}\left(z-v_{c_{n}}\right)\left(x_{c_{n}}\right)=\lim _{c_{n} \rightarrow 0}(z-v)\left(x_{c_{n}}\right)=(z-v)\left(x_{0}\right) \geq 0
$$

i.e., $x_{0} \in E(v) \subset E^{\prime}(v)$. This is a contradiction to $x_{0} \in T \backslash E^{\prime}(v)$ and shows that the minimum $m$ in (A.3) cannot be positive.


Fig. A.1. The geometric situation in the proof of Lemma A.3.
The intuition behind the proof of Lemma A. 3 is clear: If $v \in V$ is a function with $E(v) \subseteq \operatorname{conv}\left(p_{1}, \ldots, p_{k-1}, p_{k+1}, \ldots, p_{d+1}\right)$ such that there is no $\zeta \in E(v)$ with (A.1), then we can increase the value $v\left(p_{k}\right)$ without violating the constraint $0 \leq v \leq z$ and $v$ cannot be a maximal element. Using Lemma A.3, we obtain:
Theorem A.4. Let $T \subset \mathbb{R}^{d}, p_{1}, \ldots, p_{d+1}$ and $0 \leq z \in C^{1, \gamma}(T)$ be as before. Then there exists an affine linear function $v: T \rightarrow \mathbb{R}$ such that $0 \leq v \leq z$ and

$$
z\left(p_{k}\right)-v\left(p_{k}\right) \leq \frac{\sqrt{d}}{1+\gamma} \operatorname{diam}(T)^{1+\gamma}|z|_{C^{1, \gamma}(T)} \quad \forall k=1, . ., d+1
$$

Proof. If $v$ is an arbitrary maximal element of $V$ and $p_{k}$ a vertex of $T$, then there are three possibilities: If $p_{k} \in E(v)$, then it holds $v\left(p_{k}\right)=z\left(p_{k}\right)$ and the claim is certainly true. If, on the other hand, $p_{k} \notin E(v)$ and there exist $\zeta \in E(v)$ and $\varepsilon>0$ such that $\zeta+\epsilon\left(p_{k}-\zeta\right), \zeta-\epsilon\left(p_{k}-\zeta\right) \in T$, then (A.2) implies

$$
\nabla(z-v)(\zeta) \cdot\left(p_{k}-\zeta\right)=0
$$

and we may compute

$$
\begin{align*}
(z-v)\left(p_{k}\right) & =(z-v)\left(p_{k}\right)-(z-v)(\zeta) \\
& =\int_{0}^{1}\left[\nabla(z-v)\left(\zeta+t\left(p_{k}-\zeta\right)\right)-\nabla(z-v)(\zeta)\right] \cdot\left(p_{k}-\zeta\right) \mathrm{d} t \\
& =\int_{0}^{1}\left[\nabla z\left(\zeta+t\left(p_{k}-\zeta\right)\right)-\nabla z(\zeta)\right] \cdot\left(p_{k}-\zeta\right) \mathrm{d} t \\
& \leq \int_{0}^{1} \frac{\left\|\nabla z\left(\zeta+t\left(p_{k}-\zeta\right)\right)-\nabla z(\zeta)\right\|}{\left\|t\left(p_{k}-\zeta\right)\right\|^{\gamma}}\left\|p_{k}-\zeta\right\|^{\gamma+1} t^{\gamma} \mathrm{d} t \\
& \leq|z|_{C^{1, \gamma}(T)} \sqrt{d} \int_{0}^{1}\left\|p_{k}-\zeta\right\|^{\gamma+1} t^{\gamma} \mathrm{d} t \\
& \leq \frac{\sqrt{d}}{1+\gamma} \operatorname{diam}(T)^{1+\gamma}|z|_{C^{1, \gamma}(T)} \tag{A.4}
\end{align*}
$$

This proves the claim in the second case. If, lastly, $p_{k} \notin E(v)$ and there are no $\zeta \in E(v)$ and $\varepsilon>0$ such that $\zeta+\epsilon\left(p_{k}-\zeta\right), \zeta-\epsilon\left(p_{k}-\zeta\right) \in T$, then it necessarily holds $E(v) \subseteq \operatorname{conv}\left(p_{1}, \ldots, p_{k-1}, p_{k+1} \ldots, p_{n+1}\right)$ and we may employ Lemma A. 3 to obtain that there is some $\zeta \in E(v)$ with

$$
\nabla(z-v)(\zeta) \cdot\left(p_{k}-\zeta\right)=0
$$

A calculation analogous to (A.4) now yields the claim. This completes the proof.
It should be noted that there is no straightforward way to construct the approximation $v$ appearing in the last theorem from the Lagrange interpolant of the function $z$ since it is in general unclear how the interpolant has to be modified such that both the constraints $v \leq z$ and $v \geq 0$ are satisfied. We conclude our investigation with the following result about global approximations:
Corollary A.5. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded polyhedric domain with a quasi-uniform family of triangulations $\left\{\mathcal{T}_{h}\right\}$ and let $\rho_{1}$ and $\rho_{2}$ be defined as in Definition 2.2. Assume that $z \in W^{2, q}(\Omega) \cap H_{0}^{1}(\Omega)$ for some $d<q<\infty$ and suppose that a family of functions $\left\{w_{h}\right\}$ is given such that

$$
z \leq w_{h} \text { a.e. in } \Omega \quad \text { and } \quad w_{h} \in V_{h}^{0} \quad \forall h>0,
$$

where $V_{h}^{0}:=\left\{v \in C(\bar{\Omega}):\left.v\right|_{T}\right.$ is affine for all $T \in \mathcal{T}_{h}$ and $\left.\left.v\right|_{\partial \Omega}=0\right\}$. Then there exists a family of approximations $\left\{z_{h}\right\}$ satisfying

$$
z \leq z_{h} \leq w_{h} \text { a.e. in } \Omega \quad \text { and } \quad z_{h} \in V_{h}^{0}
$$

for all $h$ such that

$$
\left\|z-z_{h}\right\|_{L^{\infty}} \leq C h^{2-d / q}\|z\|_{W^{2, q}}
$$

holds with a constant $C$ independent of $h$.
Proof. Let $h$ be arbitrary but fixed, then it follows from $W^{2, q}(\Omega) \hookrightarrow C^{1,1-d / q}(\bar{\Omega})$ and Theorem A. 4 that for every $T \in \mathcal{T}_{h}$ there exists an affine linear $v_{T}: T \rightarrow \mathbb{R}$ such that $0 \leq v_{T} \leq w_{h}-z$ and

$$
0 \leq\left(w_{h}-z\right)\left(p_{k}\right)-v_{T}\left(p_{k}\right) \leq \frac{\sqrt{d}}{2-d / q} \operatorname{diam}(T)^{2-d / q}|z|_{C^{1,1-d / q}(T)}
$$

holds for all vertices $p_{k}$ of $T$. We now define $v_{h}$ to be the unique element of $V_{h}^{0}$ with

$$
v_{h}\left(x_{i}\right)=\min _{T \in \mathcal{T}_{h}: x_{i} \in T} v_{T}\left(x_{i}\right)
$$

for all mesh nodes $x_{i}$. This $v_{h}$ certainly satisfies $0 \leq v_{h} \leq v_{T} \leq w_{h}-z$ on every mesh cell $T$ and

$$
\begin{aligned}
0 \leq\left(w_{h}-z\right)\left(x_{i}\right)-v_{h}\left(x_{i}\right) & \leq \frac{\sqrt{d}}{2-d / q}\left(\max _{T \in \mathcal{T}_{h}: x_{i} \in T} \operatorname{diam}(T)^{2-d / q}|z|_{C^{1,1-d / q}(T)}\right) \\
& \leq \frac{\sqrt{d}}{2-d / q} \rho_{1}^{2-d / q} h^{2-d / q} \max _{T \in \mathcal{T}_{h}}|z|_{C^{1,1-d / q}(T)}
\end{aligned}
$$

for all nodes $x_{i}$ of the mesh. Defining $z_{h}:=w_{h}-v_{h}$, we now obtain a function with $z_{h} \in V_{h}^{0}, z \leq z_{h} \leq w_{h}$ and

$$
\left\|I_{h} z-z_{h}\right\|_{L^{\infty}} \leq \max _{x_{i}}\left|z\left(x_{i}\right)-z_{h}\left(x_{i}\right)\right| \leq C h^{2-d / q}|z|_{C^{1,1-d / q}(\bar{\Omega})}
$$

for some $C=C\left(d, q, \rho_{1}\right)>0$. The claim now follows from the triangle inequality, the Sobolev embedding $W^{2, q}(\Omega) \hookrightarrow C^{1,1-d / q}(\bar{\Omega})$ and well-known error estimates for the $L^{\infty}$-error of the Lagrange interpolant ([3, Theorem 4.4.20]).

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