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On upper bounds for regularity indices related to approximation theory
P. A. Cioica-Licht and M. Weimar

# On upper bounds for regularity indices related to approximation theory 

Petru A. Cioica-Licht* Markus Weimar ${ }^{\dagger}$

Dedicated to Prof. Dr. Stephan Dahlke on the occasion of his 60th birthday


#### Abstract

We study the interrelation between the limit $L_{p}(\Omega)$-Sobolev regularity $\bar{s}_{p}$ of (classes of) functions on bounded Lipschitz domains $\Omega \subseteq \mathbb{R}^{d}, d \geq 2$, and the limit regularity $\bar{\alpha}_{p}$ within the corresponding adaptivity scale of Besov spaces $B_{\tau, \tau}^{\alpha}(\Omega)$, where $1 / \tau=\alpha / d+1 / p$ and $\alpha>0$ ( $p>1$ fixed). The former determines the convergence rate of uniform numerical methods, whereas the latter corresponds to the convergence rate of best $N$-term approximation. We show how additional information on the Besov or TriebelLizorkin regularity may be used to deduce upper bounds for $\bar{\alpha}_{p}$ in terms of $\bar{s}_{p}$ simply by means of classical embeddings and the extension of complex interpolation to suitable classes of quasi-Banach spaces due to Kalton, Mayboroda, and Mitrea (Contemp. Math. 445). The results are applied to the Poisson equation, to the $p$-Poisson problem, and to the inhomogeneous stationary Stokes problem. In particular, we show that already established results on the Besov regularity for the Poisson equation are sharp. Keywords: Non-linear approximation, adaptive methods, Besov space, Triebel-Lizorkin space, regularity of solutions, stationary Stokes equation, Poisson equation, $p$-Poisson equation, Lipschitz domain. 2010 Mathematics Subject Classification: 35B35, 35J92, 41A25, 46E35, 65M99.

^[ *Universität Duisburg-Essen, Fakultät Mathematik, AG Stochastische Analysis, 45117 Essen AND University of Otago, Department of Mathmatics and Statistics, P.O. Box 56, Dunedin 9054, New Zealand. Email: petru.cioica-licht@uni-due.de ${ }^{\dagger}$ Corresponding author. Ruhr University Bochum, Faculty of Mathematics, Research Group Numerics, Universitätsstraße 150, 44801 Bochum, Germany. Email: markus.weimar@rub.de. ]


## 1 Introduction

The convergence rate of approximation methods strongly depends on the regularity of the target function. In particular, the convergence rate of the best $N$-term approximation for a function $f: \Omega \rightarrow \mathbb{R}$ on a bounded Lipschitz domain $\Omega \subseteq \mathbb{R}^{d}, d \in \mathbb{N}$, is intimately related to its regularity in the scale of Besov spaces

$$
\begin{equation*}
B_{\tau, \tau}^{\alpha}(\Omega), \quad \frac{1}{\tau}=\frac{\alpha}{d}+\frac{1}{p}, \quad \alpha>0 \tag{*}
\end{equation*}
$$

whereas the convergence of an approximation method based on uniform refinements depends on the regularity in the scale $W_{p}^{s}(\Omega), s>0$, of Sobolev spaces; here, $1<p<\infty$ is fixed and the approximation error is measured in $L_{p}(\Omega)$. Roughly speaking, if (and only if) the Besov regularity of the target function in the scale $(*)$ is strictly higher than its corresponding Sobolev regularity, a higher convergence rate may be achieved by switching from uniform refinement strategies to more sophisticated adaptive wavelet or finite element schemes. We refer to $[6,13,18]$ and to the references therein for details and sufficient assumptions for such statements. Definitions of the relevant function spaces are provided in the appendix.

The Sobolev regularity of solutions to elliptic partial differential equations on non-smooth domains may be very limited, even if the forcing terms are infinitely smooth. Upper bounds for

$$
\begin{equation*}
\bar{s}_{p}:=\bar{s}_{p}(S(\Omega)):=\sup \left\{s>0 \mid S(\Omega) \subseteq W_{p}^{s}(\Omega)\right\} \tag{1}
\end{equation*}
$$

where $S(\Omega) \subseteq L_{p}(\Omega)$ is a suitably chosen set of solutions to various instances of elliptic equations, can be found, for instance, in $[4,17,19,23,27,30]$. To mention an example, there exist bounded $\mathcal{C}^{1}$ domains $\Omega \subseteq \mathbb{R}^{d}$ such that if we define $S(\Omega)$ to be the set of all solutions to the Poisson equation with zero Dirichlet boundary conditions and right hand sides $f \in \mathcal{C}^{\infty}(\bar{\Omega})$, then $\bar{s}_{p}(S(\Omega))=1+1 / p$, see Section 3.1 for details. Similar results for (stochastic) evolution equations can be found, e.g., in [20, 25]. At the same time, we know that the solution to most of the equations in the aforementioned references may have higher regularity $\alpha>\bar{s}_{p}$ in the scale (*), see, e.g., $[3,5,7,8,10,11,12,15,16,21]$. For instance, in the example above, it is known that

$$
S(\Omega) \subseteq B_{\tau, \tau}^{\alpha}(\Omega), \quad \frac{1}{\tau}=\frac{\alpha}{d}+\frac{1}{p}, \quad \text { for all } \quad 0<\alpha<\left(1+\frac{1}{p}\right) \frac{d}{d-1}
$$

see [7]. The higher Besov regularity justifies the development of adaptive numerical methods for (stochastic) partial differential equations. However, to the best of our knowledge, no upper bound at all for the regularity in the scale (*), i.e., for

$$
\begin{equation*}
\bar{\alpha}_{p}:=\bar{\alpha}_{p}(S(\Omega)):=\sup \left\{\alpha>0 \mid S(\Omega) \subseteq B_{\tau, \tau}^{\alpha}(\Omega), \quad \frac{1}{\tau}=\frac{\alpha}{d}+\frac{1}{p}\right\} \tag{2}
\end{equation*}
$$

can be found in the literature; here, $\sup \emptyset:=-\infty$. Thus, in many settings, we do know that there is the possibility to outperform uniform methods by adaptive refinement strategies but we do not know how high the convergence rate of these methods can maximally get.

In this paper we study the interrelation between the limit regularity indices $\bar{s}_{p}$ and $\bar{\alpha}_{p}$. In Section 2 we prove an abstract result showing for arbitrary sets $S(\Omega) \subseteq L_{p}(\Omega)$ how additional information about the Besov or Triebel-Lizorkin regularity of all $u \in S(\Omega)$ can be used to deduce upper bounds for $\bar{\alpha}_{p}$ in terms of $\bar{s}_{p}$ simply by means of the extension of complex interpolation to suitable classes of quasi-Banach spaces from [24] and classical embeddings. We apply this result in Section 3 to the Poisson equation, the $p$-Poisson problem, and the inhomogeneous stationary Stokes equation. In particular, we show that under fairly natural assumptions, already established positive results on the Besov regularity of the solution to the Poisson equation in the scale $(*)$ are actually sharp. Before we start, we introduce some notation and comment on so-called DeVore-Triebel diagrams, which we will use in order to visualize results.

Notation. Throughout this manuscript, $\Omega$ denotes a bounded Lipschitz domain in $\mathbb{R}^{d}$ for some $d \in \mathbb{N}$. For $0<p<\infty$, by $L_{p}(\Omega)$ we denote the space of all (equivalence classes of) Lebesgue-measurable, scalar-valued functions satisfying $\left\|\left.u\left|L_{p}(\Omega) \|^{p}:=\int_{\Omega}\right| u(x)\right|^{p} \mathrm{~d} x<\infty\right.$, while $L_{\infty}(\Omega)$ is the space of all (equivalence classes of) Lebesgue-measurable, Lebesguealmost everywhere bounded scalar-valued functions on $\Omega$. Moreover, $B_{p, q}^{s}(\Omega)$ and $F_{p, q}^{s}(\Omega)$ stand for the Besov and Triebel-Lizorkin spaces, respectively, with smoothness parameter $s \in \mathbb{R}$, integrability parameter $p \in(0, \infty]$ (with $p<\infty$ for Triebel-Lizorkin spaces) and microscopic parameter $q \in(0, \infty]$. The corresponding spaces $B_{p, q}^{s}(\partial \Omega)$ and $F_{p, q}^{s}(\partial \Omega)$ on the boundary $\partial \Omega$ of the domain $\Omega$ are defined as in [28]. For $1<p<\infty$, by $W_{p}^{s}(\Omega)$ we denote the $L_{p}(\Omega)$-Sobolev space of order $s \in \mathbb{R}$. For two quasi-normed spaces $X$ and $Y$, we write $X \hookrightarrow Y$ if $X$ is continuously and linearly embedded in $Y$ and $[X, Y]_{\theta}$ stands for the complex interpolation space of the pair $(X, Y)$ with parameter $\theta \in(0,1)$. Precise definitions and relevant interpolation and embedding properties of Besov, Triebel-Lizorkin, and Sobolev spaces are collected in Appendix A.

Throughout, the letter $C$ is used to denote a finite positive constant that may differ from one appearance to another, even in the same chain of inequalities. Moreover, we adopt the usual conventions $1 / \infty:=0$ and $1 / 0:=\infty$.

DeVore-Triebel diagrams. We are going to use so-called DeVore-Triebel diagrams in order to visualize results. In those $(1 / p, s)$-diagrams, we identify every point $(1 / p, s) \in[0, \infty) \times \mathbb{R}$ with the Besov space $B_{p, p}^{s}(\Omega)$. Many embedding and interpolation results for Besov spaces can then be visualized in a very convenient way (see Figure 1):


Figure 1: Visualization of Besov spaces on bounded Lipschitz domains $\Omega \subseteq \mathbb{R}^{d}$ in a DeVore-Triebel diagram.

- Besov spaces form scales of (generalized) complex interpolation spaces, see Proposition A.4. As a consequence, if $f \in B_{p_{i}, p_{i}}^{s_{i}}(\Omega)$ for $i=0,1$, then $f \in B_{\widetilde{p}, \widetilde{p}}^{\widetilde{s}}(\Omega)$ for all $(1 / \widetilde{p}, \widetilde{s})$ on the line segment between $\left(1 / p_{0}, s_{0}\right)$ and $\left(1 / p_{1}, s_{1}\right)$; see (i) in Figure 1.
- If $f \in B_{p, p}^{s}(\Omega)$ for some $0<p<\infty$ and $s \in \mathbb{R}$, then, by Proposition A.3(iv), $f$ is contained in all the Besov spaces represented by the points $(1 / \widetilde{p}, \widetilde{s}) \in[0, \infty) \times \mathbb{R}$ with $\widetilde{s}<s-d \max \{1 / p-1 / \widetilde{p}, 0\}$; see the shaded area (ii) in Figure 1. Moreover, by Proposition A.3(v), it is contained in all Besov spaces represented by the points $(1 / \widetilde{p}, \widetilde{s}) \in(0,1 / p) \times \mathbb{R}$ with $\widetilde{s}=s-d(1 / p-1 / \widetilde{p})$; see (iii) in Figure 1.
- If $f \in A_{p_{z}, q_{z}}^{z}(\Omega)$ for some $A \in\{B, F\}, z \in \mathbb{R}$ and $0<p_{z}, q_{z} \leq \infty$ (with finite $p_{z}$ if $A=F$ ), then, by Proposition A.3(iv), $f$ is contained in all Besov spaces represented by the ray $\left\{\left(1 / p_{z}, \widetilde{s}\right) \mid \widetilde{s}<z\right\}$; see (iv) in Figure 1.
Moreover, in such a diagram, for $1<p<\infty$, the scale $(*)$ is represented by the so-called $L_{p}(\Omega)$-Sobolev embedding line

$$
\begin{equation*}
\left\{\left(\frac{1}{\tau}, \alpha\right) \in(0, \infty)^{2} \left\lvert\, \frac{1}{\tau}=\frac{\alpha}{d}+\frac{1}{p}\right.\right\} \tag{3}
\end{equation*}
$$

see (v) in Figure 1.

## 2 Main result

In this section we analyze how additional information about the Besov or Triebel-Lizorkin regularity may be used in order to derive upper bounds for $\bar{\alpha}_{p}$ in terms of $\bar{s}_{p}$ simply by means of complex interpolation and classical embedding theorems; here and in the sequel, $\bar{s}_{p}$ and $\bar{\alpha}_{p}$ are as defined in Section 1, see (1) and (2), respectively. We prove the following main result.
Theorem 2.1. For $d \in \mathbb{N}$ let $\Omega \subseteq \mathbb{R}^{d}$ be a bounded Lipschitz domain. Moreover, let $1<p<\infty$ and let $S(\Omega) \subseteq L_{p}(\Omega)$ be such that $0<\bar{s}_{p}=\bar{s}_{p}(S(\Omega)) \leq \infty$. Assume that for some $z \in \mathbb{R}$ and some $p<p_{z} \leq \infty, S(\Omega) \subseteq B_{p_{z}, p_{z}}^{s}(\Omega)$ for all $s<z$. Then

$$
\begin{equation*}
z \leq \bar{s}_{p} \leq \bar{\alpha}_{p} . \tag{4}
\end{equation*}
$$

If additionally

$$
\mu:=\mu\left(p_{z}, p, \bar{s}_{p}, d\right):=\bar{s}_{p}-d\left(\frac{1}{p}-\frac{1}{p_{z}}\right)<z,
$$

then

$$
\begin{equation*}
\bar{\alpha}_{p} \leq \bar{s}_{p}+\bar{s}_{p} \cdot \frac{\bar{s}_{p}-z}{z-\mu}=\bar{s}_{p} \cdot \frac{\bar{s}_{p}-\mu}{z-\mu} . \tag{5}
\end{equation*}
$$

Before we give a proof of this theorem, let us make some remarks. We start with a sufficient condition for the additional regularity assumption.
Remark 2.2. Let $0<p_{z}<\infty$ and $z \in \mathbb{R}$. Then, by classical embedding theorems for Besov and Triebel-Lizorkin spaces, as collected in Proposition A.3, the assertion

$$
S(\Omega) \subseteq A_{p_{z}, q_{z}}^{z}(\Omega) \quad \text { for some } \quad A \in\{B, F\} \quad \text { and } \quad 0<q_{z} \leq \infty
$$

is sufficient for

$$
S(\Omega) \subseteq B_{p_{z}, p_{z}}^{s}(\Omega) \quad \text { for all } \quad s<z
$$

Moreover, so is

$$
S(\Omega) \subseteq A_{p_{z}, q_{z}}^{s}(\Omega) \quad \text { for some } \quad A \in\{B, F\}, \quad 0<q_{z} \leq \infty, \quad \text { and all } \quad s<z
$$

If $A=B$, then these implications also hold for $p_{z}=\infty$.
Remark 2.3. In principle, $S(\Omega)$ could be any subset of some Besov/Triebel-Lizorkin space. But even if we restrict ourselves to solution sets for operator equations, there are several different interpretations: On the one hand, we may think of one particular problem given by a fixed operator $L$ acting on functions defined on a fixed domain $\Omega$ with fixed right-hand side and fixed initial/boundary conditions if necessary. Then $S(\Omega)$ only contains solutions for this particular situation and we probably even have $\# S(\Omega)=1$ such that $\bar{s}_{p}$ and $\bar{\alpha}_{p}$ describe smoothness properties of one particular function. On the other hand, we may also think of solution sets for classes of problems such as, e.g.,
(i) a fixed equation (like the Poisson equation $\Delta u=f$ with zero Dirichlet boundary condition $u_{\mid \partial \Omega}=0$ ) on a fixed domain $\Omega$ (e.g., the standard L-shape domain in $d=2$ ) with variable right-hand side from a certain class of functions (e.g., arbitrary $f \in L_{2}(\Omega)$ ), or
(ii) a class of operator equations (e.g., all linear, second order PDEs with smooth coefficients) on a fixed domain $\Omega$ with, say, smooth right-hand sides,
and so forth. Since in this case $S(\Omega)$ collects all functions which solve at least one admissible problem instance, here $\bar{s}_{p}$ and $\bar{\alpha}_{p}$ describe lower bounds for the regularity of solutions to the hardest possible problem in the respective class. For example, $u^{*} \equiv 0$ solves the problem described in (i) for $f \equiv 0$. Hence, $u^{*} \in S(\Omega)$ and $u^{*} \in \bigcap_{s>0} W_{2}^{s}(\Omega)$, but $\bar{s}_{2}=5 / 3<\infty$, see also Remark 3.3 below.

We could even go one step further and consider classes of problems like
(iii) a fixed equation considered on a class of domains (e.g., all bounded $\mathcal{C}^{1}$ domains) with certain restrictions on the right-hand side and/or on initial/boundary conditions.
However, then the notation would get more complicated such that in the sequel we restrict ourselves to the cases mentioned above.

Remark 2.4. Throughout this remark, we assume that we are in the setting of Theorem 2.1.
(i) Note that, due to standard embeddings of Besov and Triebel-Lizorkin spaces (as provided in Proposition A. 3 in the appendix), for $A \in\{B, F\}$ and $0<q \leq \infty$ we have

$$
\bar{s}_{p}=\bar{s}_{p}(S(\Omega))=\bar{s}_{p, q, A}(S(\Omega)):=\sup \left\{s>0 \mid S(\Omega) \subseteq A_{p, q}^{s}(\Omega)\right\} .
$$

That is, the limit regularity index $\bar{s}_{p, q, A}$ does not depend on the microscopic parameter $q$, nor on the type of the spaces $A$ (Besov vs. Triebel-Lizorkin). Moreover, it coincides with $\bar{s}_{p}$ defined in (1). In particular,

$$
\begin{equation*}
\bar{s}_{p}=\bar{s}_{p, p, B}(\Omega)=\sup \left\{s>0 \mid S(\Omega) \subseteq B_{p, p}^{s}(\Omega)\right\} \tag{6}
\end{equation*}
$$

and also

$$
\bar{s}_{p}=\bar{s}_{p, \infty, B}(S(\Omega))=\sup \left\{s>0 \mid S(\Omega) \subseteq B_{p, \infty}^{s}(\Omega)\right\}
$$

where the latter quantity is defined by means of the slightly larger Besov spaces $B_{p, \infty}^{s}(\Omega)$ which coincide with the approximation spaces $\mathcal{A}_{\infty}^{s / d}\left(L_{p}(\Omega)\right)$ w.r.t. non-adaptive algorithms based on uniform refinement, see, e.g., [13] for details.
(ii) Due to the generalization of Sobolev's embedding theorem to Besov spaces (as presented in Proposition A.3(v)), a space $B_{\tau, \tau}^{\alpha}(\Omega)$ from the adaptivity scale $(*)$ is embedded into every other space $B_{\tau_{0}, \tau_{0}}^{\alpha_{0}}(\Omega), 1 / \tau_{0}=\alpha_{0} / d+1 / p$, from the same scale ( $*$ ) with $\alpha_{0}<\alpha$.


Figure 2: Visualization of statement and proof of Assertion (5) from Theorem 2.1 in a DeVore-Triebel diagram.

However, as a consequence of the sharpness of Sobolev embeddings, the space $B_{\tau, \tau}^{\alpha}(\Omega)$ is not embedded in $A_{p, q}^{s}(\Omega)$ for any $A \in\{B, F\}, 0<q \leq \infty$, and $s>0$, as this would contradict the 'only if' part of Proposition A.3(v). As a consequence, it is not possible to obtain a non-trivial upper bound for $\bar{\alpha}_{p}$ in terms of $\bar{s}_{p}$ without further assumptions on $S(\Omega)$.
(iii) In Figure 2 we use a DeVore-Triebel diagram to visualize our upper bound (5) for $\bar{\alpha}_{p}$ and the corresponding proof idea, given that $\bar{s}_{p}<\infty$; relevant details on DeVoreTriebel diagrams can be found at the end of the introduction above. The bound $\bar{s}_{p} \cdot\left(\bar{s}_{p}-\mu\right) /(z-\mu)$ in (5) is precisely the ordinate of the intersection point of the (dashed) line through $\left(1 / p_{z}, z\right)$ and $\left(1 / p, \bar{s}_{p}\right)$ with the $L_{p}(\Omega)$-Sobolev embedding line (3). Therefore, by elementary geometry, for every $\alpha>\bar{s}_{p} \cdot\left(\bar{s}_{p}-\mu\right) /(z-\mu)$, there exists $\widetilde{z}<z$, such that the (solid) line through $\left(1 / p_{z}, \widetilde{z}\right)$ and $(\alpha / d+1 / p, \alpha)$ contains a point $(1 / p, s)$ for some $s>\bar{s}_{p}$. Since $S(\Omega) \subseteq B_{p_{z}, p_{z}}^{\tilde{z}}(\Omega)$ for all $\widetilde{z}<z$, the claim $S(\Omega) \subseteq B_{\tau, \tau}^{\alpha}(\Omega)$ for such an $\alpha$ would thus contradict the maximality of $\bar{s}_{p}$, see also (6).
(iv) The proof idea above obviously fails if $z \leq \mu\left(p_{z}, p, \bar{s}_{p}, d\right)$, i.e., if the point $\left(1 / p_{z}, z\right)$ is below or exactly on the Sobolev embedding line $\left\{(1 / \widetilde{p}, \widetilde{s}) \mid \widetilde{s}=\bar{s}_{p}-d(1 / p-1 / \widetilde{p})\right\}$ through $\left(1 / p, \bar{s}_{p}\right)$. In this case the line through $\left(1 / p_{z}, z\right)$ and $\left(1 / p, \bar{s}_{p}\right)$ does not intersect with the corresponding $L_{p}(\Omega)$-Sobolev embedding line (3).

Actually, it is clear that we cannot even expect to obtain a non-trivial bound on $\bar{\alpha}_{p}$ if we only know that $S(\Omega) \subseteq B_{p_{z}, p_{z}}^{s}(\Omega)$ for all $s<z \leq \mu\left(p_{z}, p, \bar{s}_{p}, d\right)$, since this is already implied by Sobolev's embedding theorem (see Proposition A.3(iv)). Thus, assuming this does not add any additional information about $S(\Omega)$ and we cannot expect to be able to establish a non-trivial bound on $\bar{\alpha}_{p}$, see also (ii) above. In the limiting case, i.e., if $z=\mu$, then assuming that $S(\Omega) \subseteq A_{p_{z}, q_{z}}^{z}(\Omega)$ for some $A \in\{B, F\}$ and $0<q_{z} \leq \infty$ as in Remark 2.2 may or may not constitute an additional assumption on $S(\Omega)$. However, also in this case it is not possible to establish a non-trivial bound for $\bar{\alpha}_{p}$. Counterexamples can easily be constructed in terms of standard representatives of Besov and Triebel-Lizorkin spaces; see, in particular, [29, Lemma 2.3.1.1].
(v) The proof technique described in (iii) above may also be used in order to derive, for instance,

- the lower bound

$$
\widetilde{s}_{p}:=\bar{\alpha}_{p} \cdot \frac{z+d\left(1 / p-1 / p_{z}\right)}{\bar{\alpha}_{p}+d\left(1 / p-1 / p_{z}\right)}
$$

for $\bar{s}_{p}$, provided we are given $\bar{\alpha}_{p}>0$ and $S(\Omega) \subseteq A_{p_{z}, q_{z}}^{z}(\Omega)$ for some $A \in\{B, F\}$, $p_{z}>p, 0<q_{z} \leq \infty$, and $z \in \mathbb{R}$, or

- an upper bound for $\bar{s}_{\widehat{p}}$ for some $\widehat{p}>p$, given $\bar{s}_{p}$, as well as $S(\Omega) \subseteq A_{p_{z}, q_{z}}^{z}(\Omega)$ for some $A \in\{B, F\}, z>\bar{s}_{p}$, and $p_{z}<p$.

In Section 3.1, we are going to use the latter in order to determine $\bar{s}_{p}, 1<p<\infty$, for the Poisson equation with smooth right-hand sides and zero Dirichlet boundary conditions on a bounded $\mathcal{C}^{1}$ domain constructed by Jerison and Kenig [23].
(vi) Further assumptions of the type $S(\Omega) \subseteq A_{\tilde{\tilde{p}_{z}}, \widetilde{q}_{F}}^{\tilde{q}_{F}}(\Omega)$ for some $A \in\{B, F\}$, as well as $1<p<p_{z}<\widetilde{p}_{z} \leq \infty$ (with finite $\widetilde{p}_{z}$ if $A=F$ ), $0<\widetilde{q}_{z} \leq \infty$, and $\widetilde{z} \in \mathbb{R}$ lead to an improvement of the upper bound for $\bar{\alpha}_{p}$ by means of the proof technique described in (iii) only if the point $\left(1 / \widetilde{p}_{z}, \widetilde{z}\right)$ lies strictly above the line through the two points $\left(1 / p_{z}, z\right)$ and $\left(1 / p, \bar{s}_{p}\right)$ in the DeVore-Triebel diagram. Moreover, by complex interpolation it becomes obvious that the set of parameters

$$
\left\{\left.\left(\frac{1}{\varrho}, s\right) \in[0, \infty)^{2} \right\rvert\, S(\Omega) \subseteq B_{\varrho, \varrho}^{s}(\Omega)\right\}
$$

is necessarily convex and that each $\left(1 / \varrho, \bar{s}_{\varrho}\right)$ with $0<\varrho \leq \infty$ belongs to its boundary.
(vii) For $1<p<\infty$, the regularity of a function in the scale $(*)$ is intimately related to the convergence rate of the best $N$-term approximation, if the error is measured in $L_{p}(\Omega)$.

However, if the error is to be measured in the norm of some other Sobolev space $W_{p}^{r}(\Omega)$ with $r>0$ (describing, for instance, the energy space), then the scale changes to

$$
B_{\tau, \tau}^{\alpha}(\Omega), \quad \frac{1}{\tau}=\frac{\alpha-r}{d}+\frac{1}{p}, \quad \alpha>r .
$$

Since this is just a shift of the $L_{p}(\Omega)$-Sobolev embedding line, our analysis carries over to this case mutatis mutandis. For the ease of presentation we omit the details. Moreover, we can replace the underlying Lipschitz domain $\Omega$ by a (patchwise smooth) manifold; cf. [9, 12, 34].

Now we are ready to give a detailed proof of Theorem 2.1.
Proof of Theorem 2.1. Relation (4) follows by contradiction due to the fact that for all $0<p_{1}<p_{0}<\infty$ and $s_{1}<s_{0}$ there holds $B_{p_{0}, p_{0}}^{s_{0}}(\Omega) \hookrightarrow B_{p_{1}, p_{1}}^{s_{1}}(\Omega)$, see Proposition A.3(iv). This embedding also implies that $\bar{\alpha}_{p}=\infty$ if $\bar{s}_{p}=\infty$. Thus, we are left with proving (5) for $\bar{s}_{p}<\infty$. Again we argue by contradiction. Assume $S(\Omega) \subseteq B_{\tau, \tau}^{\alpha}(\Omega), 1 / \tau=\alpha / d+1 / p$, for some $\alpha>\bar{s}_{p} \cdot\left(\bar{s}_{p}-\mu\right) /(z-\mu)$. Since $S(\Omega) \subseteq B_{p_{z}, p_{z}}^{\tilde{z}}(\Omega)$ for all $\widetilde{z}<z$, we also know that $S(\Omega) \subseteq B_{\widetilde{p}, \tilde{p}}^{\widetilde{s}}(\Omega)$ with $\widetilde{s}=(1-\theta) \widetilde{z}+\theta \alpha$ and $1 / \widetilde{p}=(1-\theta) / p_{z}+\theta / \tau$ for all $\theta \in(0,1)$, see Proposition A.4. In particular, if we choose

$$
\theta=\theta_{0}:=\frac{1 / p-1 / p_{z}}{1 / \tau-1 / p_{z}}=\frac{1 / p-1 / p_{z}}{\alpha / d+1 / p-1 / p_{z}}=\frac{\bar{s}_{p}-\mu}{\alpha+\bar{s}_{p}-\mu} \in(0,1)
$$

we obtain $S(\Omega) \subseteq B_{p, p}^{\widetilde{s}}(\Omega)$ for all $\widetilde{s}=\left(1-\theta_{0}\right) \widetilde{z}+\theta_{0} \alpha$ with $\widetilde{z}<z$. Since $\alpha>\bar{s}_{p} \cdot\left(\bar{s}_{p}-\mu\right) /(z-\mu)$, we have

$$
\left(1-\theta_{0}\right) z+\theta_{0} \alpha=\frac{\alpha\left(z+\bar{s}_{p}-\mu\right)}{\alpha+\bar{s}_{p}-\mu}=\frac{z+\bar{s}_{p}-\mu}{1+\left(\bar{s}_{p}-\mu\right) / \alpha}>\frac{z+\bar{s}_{p}-\mu}{\left(z+\bar{s}_{p}-\mu\right) / \bar{s}_{p}}=\bar{s}_{p} .
$$

Therefore, there exists $\widetilde{z}<z$, such that $s:=\left(1-\theta_{0}\right) \widetilde{z}+\theta_{0} \alpha>\bar{s}_{p}$, which means that $S(\Omega) \subseteq B_{p, p}^{s}(\Omega)$ for some $s>\bar{s}_{p}$. But this contradicts the maximality of $\bar{s}_{p}$.

## 3 Examples

In this section we apply Theorem 2.1 to three sample problems: the Poisson equation, the $p$-Poisson problem, and the inhomogeneous stationary Stokes equation.

### 3.1 The Poisson equation

Let us consider the Poisson equation with zero Dirichlet boundary conditions

$$
\left.\begin{array}{rl}
\Delta u=f & \text { on } \Omega  \tag{7}\\
u=0 & \text { on } \partial \Omega,
\end{array}\right\}
$$

on a bounded Lipschitz domain $\Omega \subseteq \mathbb{R}^{d}, d \geq 2$. Singularities of the boundary $\partial \Omega$ of the underlying domain $\Omega$ are known to have negative effects on the regularity of the solution $u$ to (7). While on smooth domains we have the usual shift

$$
f \in W_{p}^{s-2}(\Omega) \quad \Longrightarrow \quad u \in W_{p}^{s}(\Omega),
$$

this mechanism fails if we allow the boundary of $\Omega$ to be merely $\mathcal{C}^{1}$. In this case, for instance, $f \in W_{2}^{-1 / 2}(\Omega)$ does not necessarily imply $u \in W_{2}^{3 / 2}(\Omega)$. This problem has been intensively studied in [23] by Jerison and Kenig; see also [17, 26]. Therein one may find a precise description of the range of parameters $(1 / p, s)$ that allow for shift theorems for Equation (7) in Bessel potential spaces and in Besov spaces. The sharpness of this range is underpinned by several counterexamples, see, in particular, [23, Section 6]. Motivated by these results and by the relevance of the regularity in Sobolev spaces and in the scales $(*)$ of Besov spaces in (non-)linear approximation theory, Dahlke and DeVore [7] analyzed the regularity of the Poisson equation in Besov spaces with integrability parameter less than one. Put together, the positive results from [23] and [7] guarantee the following: If we are only interested in the effect of the boundary singularities and therefore assume that $f \in \mathcal{C}^{\infty}(\bar{\Omega})$, then the solution $u \in W_{2,0}^{1}(\Omega)$ to the corresponding equation (7) is contained in every Besov space $B_{q, q}^{r}(\Omega)$ represented by a point $(1 / q, r)$ within the shaded area in the DeVore-Triebel diagram in Figure 3. Using the terminology from the previous sections, we set

$$
\begin{equation*}
S(\Omega):=\left\{u \in W_{2,0}^{1}(\Omega) \mid \Delta u \in \mathcal{C}^{\infty}(\bar{\Omega})\right\} \tag{8}
\end{equation*}
$$

Then

$$
S(\Omega) \subseteq B_{q, q}^{r}(\Omega) \quad \text { for all } \quad 0<r<1+\frac{1}{q} \quad \text { and } \quad 0<\frac{1}{q}<\frac{d+1}{d-1}
$$

such that, in particular,

$$
\begin{equation*}
\bar{s}_{p}(S(\Omega)) \geq 1+\frac{1}{p} \quad \text { and } \quad \bar{\alpha}_{p}(S(\Omega)) \geq\left(1+\frac{1}{p}\right) \frac{d}{d-1} \tag{9}
\end{equation*}
$$

for every $1<p<\infty$. The following theorem asserts the existence of bounded $\mathcal{C}^{1}$ domains on which these lower bounds for $\bar{s}_{p}$ and $\bar{\alpha}_{p}$ become also upper bounds.


Figure 3: Visualization of the Besov regularity of the Poisson equation with smooth right-hand side on bounded $\mathcal{C}^{1}$ domains in a DeVore-Triebel diagram.

Theorem 3.1. For $d \geq 2$, there exists a bounded $\mathcal{C}^{1}$ domain $\Omega \subseteq \mathbb{R}^{d}$ such that if $S(\Omega)$ is defined as in (8), then for arbitrary $1<p<\infty$ there holds

$$
\bar{s}_{p}(S(\Omega))=1+\frac{1}{p} \quad \text { and } \quad \bar{\alpha}_{p}(S(\Omega))=\left(1+\frac{1}{p}\right) \frac{d}{d-1} .
$$

Our proof of Theorem 3.1 below is based on a counterexample by Jerison and Kenig of a $\mathcal{C}^{1}$ domain $\Omega \subseteq \mathbb{R}^{d}, d \geq 2$, for which there exists a function $f \in \mathcal{C}^{\infty}(\bar{\Omega})$, such that the second derivatives of the solution $u \in W_{2,0}^{1}(\Omega)$ to the corresponding equation (7) are not contained in $L_{1}(\Omega)$, thus $u \notin W_{1}^{2}(\Omega)$. We refer to [23, Theorem $\left.1.2(\mathrm{~b})\right]$ for the statement and to $[23$, Section 6] for the corresponding counterexample. For such a solution to (7) we prove the following.

Lemma 3.2. Let $d \geq 2$. Moreover, let $\Omega \subseteq \mathbb{R}^{d}$ be a $\mathcal{C}^{1}$ domain for which there exists a function $f \in \mathcal{C}^{\infty}(\bar{\Omega})$ such that the unique solution $u \in W_{2,0}^{1}(\Omega)$ to the corresponding Poisson equation (7) satisfies $u \notin W_{1}^{2}(\Omega)$. Then the following statements hold.
(i) $u \notin B_{1,1}^{2}(\Omega)$.
(ii) If $1<p<\infty$ and $s>1+\frac{1}{p}$, then $u \notin B_{p, p}^{s}(\Omega)$.
(iii) $u \in F_{p, 2}^{1+1 / p}(\Omega)$ for all $2 \leq p<\infty$.
(iv) If $1<p<\infty$ and $\alpha \geq\left(1+\frac{1}{p}\right) \frac{d}{d-1}$, then $u \notin B_{\tau, \tau}^{\alpha}(\Omega), \frac{1}{\tau}=\frac{\alpha}{d}+\frac{1}{p}$.

Proof. We prove the four statements successively.
(i). The assertion $u \in B_{1,1}^{2}(\Omega)$ would contradict our assumption that $u \notin W_{1}^{2}(\Omega)$ since $B_{1,1}^{2}(\Omega) \hookrightarrow W_{1}^{2}(\Omega)$, which follows, e.g., from [31, Theorem 2.3.8(i) \& Proposition 2.5.7(i)].
(ii). Suppose that $u \in B_{p, p}^{s}(\Omega)$ for some $1<p<\infty$ and $s>1+1 / p$. W.l.o.g. we may also assume that $s<2$. From [7, Theorem 4.1] we can deduce that $u \in B_{q, q}^{r}(\Omega)$ with $1 / q=1+\varepsilon$ and $r=2+\varepsilon(2-s) /(1-1 / q)$ for all $0<\varepsilon<2 /(d-1)$. Then by Proposition A. 4 we have

$$
u \in\left[B_{p, p}^{s}(\Omega), B_{q, q}^{r}(\Omega)\right]_{\theta}=B_{1,1}^{2}(\Omega) \quad \text { for } \quad \theta=\frac{1-1 / p}{1-1 / p+\varepsilon} \in(0,1)
$$

However, this contradicts (i).
(iii). We prove this assertion with an argument used in [4, page 3]: Let us extend $f$ to the whole of $\mathbb{R}^{d}$ such that the extension (also denoted by $f$ ) is at least smooth enough to be contained in $F_{p, 2}^{-1+1 / p+\varepsilon}\left(\mathbb{R}^{d}\right)$ for some $\varepsilon>0$. Then the equation $\Delta v=f$ on $\mathbb{R}^{d}$ has a unique solution $v \in F_{p, 2}^{1+1 / p+\varepsilon}\left(\mathbb{R}^{d}\right)$ and $\left.v\right|_{\partial \Omega} \in B_{p, p}^{1+\varepsilon}(\partial \Omega) \hookrightarrow F_{p, 2}^{1}(\partial \Omega)$. Therefore, $\widetilde{u}:=$ $v-u$ is a harmonic function on $\Omega$ with trace $\left.\widetilde{u}\right|_{\partial \Omega} \in F_{p, 2}^{1}(\partial \Omega)$. From [23, Theorem 5.15(b)] it thus follows that $\widetilde{u} \in F_{p, 2}^{1+1 / p}(\Omega)$ and hence also $u=\widetilde{u}-v \in F_{p, 2}^{1+1 / p}(\Omega)$.
(iv). Theorems 1.1 and 1.3 of [23] together with part (ii) imply that $\bar{s}_{p}:=\bar{s}_{p}(\{u\})=1+1 / p$ for all $1<p<\infty$. Now fix $1<p<p_{z}<\infty$. Then, we may apply Theorem 2.1 with $z:=\bar{s}_{p_{z}}=1+1 / p_{z}$ and

$$
\mu=\bar{s}_{p}-d\left(\frac{1}{p}-\frac{1}{p_{z}}\right)=1+\frac{1}{p}-d\left(\frac{1}{p}-\frac{1}{p_{z}}\right)=1+\frac{1}{p_{z}}-(d-1)\left(\frac{1}{p}-\frac{1}{p_{z}}\right)<1+\frac{1}{p_{z}},
$$

to obtain

$$
\bar{\alpha}_{p}(\{u\}) \leq \bar{s}_{p} \cdot \frac{\bar{s}_{p}-\mu}{\bar{s}_{p_{z}}-\mu}=\left(1+\frac{1}{p}\right) \frac{d}{d-1}=: \widetilde{\alpha}_{p}
$$

which obviously proves (iv) if $\alpha>\widetilde{\alpha}_{p}$.

The fact that $u \notin B_{p, p}^{\widetilde{\alpha}_{p}}(\Omega)$ follows from parts (i) and (iii) by another complex interpolation argument: Since $u \in F_{2,2}^{3 / 2}(\Omega)=B_{2,2}^{3 / 2}(\Omega)$ and the points $(1 / 2,3 / 2),(1,2)$, and $\left(\widetilde{\alpha}_{p} / d+1 / p, \widetilde{\alpha}_{p}\right)$ lie on the same line of slope 1 through $(0,1)$ in a DeVore-Triebel diagram, the statement $u \in B_{p, p}^{\widetilde{\alpha}_{p}}(\Omega)$ would contradict (i).
Proof of Theorem 3.1. Due to Jerison and Kenig [23, Theorem 1.2(b)], there exist $\Omega \subseteq \mathbb{R}^{d}$ and $f \in \mathcal{C}^{\infty}(\bar{\Omega})$, such that the assumptions of Lemma 3.2 are satisfied. Therefore, the assertion follows from Lemma 3.2 and (9).

We conclude this subsection with some further remarks.
Remark 3.3. It is worth mentioning that the bounds in Theorem 3.1 are due to worst-case scenarios regarding the behaviour of $\mathcal{C}^{1}$ boundaries. However, for large classes of domains, which are not even necessarily of class $\mathcal{C}^{1}$, the regularity indices $\bar{s}_{p}(S(\Omega))$ and $\bar{\alpha}_{p}(S(\Omega))$ with $S(\Omega)$ as defined in (8) may be higher, at least for certain $1<p<\infty$. For instance, if $\Omega \subseteq \mathbb{R}^{2}$ is a polygonal domain with maximal interior angle $\kappa_{0} \in(\pi, 2 \pi)$, then Grisvard [19, 20] shows that

$$
\begin{equation*}
\bar{s}_{p}(S(\Omega))=\frac{2}{p}+\frac{\pi}{\kappa_{0}}, \quad 1<p<\infty \tag{10}
\end{equation*}
$$

which is strictly greater than $1+1 / p$ whenever $p<\kappa_{0} /\left(\kappa_{0}-\pi\right)$. Moreover, it is known from [5] that

$$
\bar{\alpha}_{2}(S(\Omega))=\infty .
$$

Note that this does not contradict Theorem 2.1 since (10) implies that for any fixed $1<p<\infty$ and all $p_{z}>p$, there is no $z>\mu\left(p_{z}, p, \bar{s}_{p}, 2\right)$ such that $S(\Omega) \subseteq B_{p_{z}, p_{z}}^{s}(\Omega)$ for all $s<z$.
Remark 3.4. In [4] Costabel constructs bounded $\mathcal{C}^{1}$ domains $\Omega \subseteq \mathbb{R}^{d}$ of arbitrary dimension $d \geq 2$, for which there exists $f \in \mathcal{C}^{\infty}(\bar{\Omega})$ such that the solution $u$ to the corresponding Poisson equation (7) is contained in $W_{2}^{3 / 2}(\Omega)$, but not in $W_{p}^{1+1 / p+\varepsilon}(\Omega)$ for any $1 \leq p<\infty$ and any $\varepsilon>0$; see, in particular, Theorem 1.2 and Remark 1.3 therein. Lemma 3.2 above shows that the counterexample provided by Jerison and Kenig in [23, Section 6] as a proof of Theorem 1.2(b) therein has these properties, too.

### 3.2 The $p$-Poisson problem

Our second example is the $p$-Poisson problem for some fixed $1<p<\infty$. For $d \geq 2$, let again $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded Lipschitz domain. Given $f \in W_{p^{\prime}}^{-1}(\Omega)$ with $1 / p+1 / p^{\prime}=1$, we
seek the unique weak solution $u \in W_{p, 0}^{1}(\Omega)$ to

$$
\left.\begin{array}{rll}
\Delta_{p} u=f & \text { on } & \Omega,  \tag{11}\\
u=0 & \text { on } & \partial \Omega,
\end{array}\right\}
$$

where $\Delta_{p} u:=\operatorname{div}\left(|\nabla u|_{2}^{p-2} \nabla u\right)$ denotes the $p$-Laplace operator.
For this problem various local and global regularity results are known; we refer, e.g., to $[1,8,14,22,30]$ and the references therein. Our subsequent analysis relies on the following result.

Proposition 3.5 (Ebmeyer [14, Theorem 2.4]). For $d \geq 2$ let $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded polyhedral Lipschitz domain. Moreover, let $1<p \leq 2$ and $f \in L_{p^{\prime}}(\Omega)$. Then the unique weak solution to (11) satisfies

$$
\begin{equation*}
u \in W_{p_{z}}^{s}(\Omega) \quad \text { for all } \quad s<\frac{3}{2} \quad \text { and } \quad p_{z}:=p_{z}(d, p):=\frac{p}{1-(2-p) /(2 d)} . \tag{12}
\end{equation*}
$$

Although, to the best of our knowledge, even in this restricted setting the exact value of $\bar{s}_{p}$ is unknown, we can apply our main Theorem 2.1 in order to deduce the following statement:

Theorem 3.6. For $d \geq 2$ let $\Omega \subseteq \mathbb{R}^{d}$ denote some bounded polyhedral Lipschitz domain. Given $1<p<2$ let $S(\Omega)$ denote the set of solutions to the p-Poisson problem (11) with right-hand sides $f \in L_{p^{\prime}}(\Omega)$. Then for the regularity indices $\bar{s}_{p}$ and $\bar{\alpha}_{p}$ as defined in (1) and (2), respectively, one of the following cases applies:
1.) $3 / 2 \leq \bar{s}_{p}<1+1 / p$ and

$$
\bar{s}_{p} \leq \bar{\alpha}_{p} \leq \bar{s}_{p} \frac{1+1 / p-3 / 2}{1+1 / p-\bar{s}_{p}} .
$$

2.) $1+1 / p \leq \bar{s}_{p} \leq \bar{\alpha}_{p}$.

Proof. For $1<p<2$ the parameter $p_{z}$ in (12) is strictly larger than $p$. Using that $W_{p_{z}}^{s}(\Omega)=B_{p_{z}, p_{z}}^{s}(\Omega)$ for $0<s \notin \mathbb{N}$, we thus can apply Theorem 2.1 with this $p_{z}$ and $z:=3 / 2$. This yields that in any case there holds

$$
\frac{3}{2} \leq \bar{s}_{p} \leq \bar{\alpha}_{p}
$$

Moreover, $\mu=\mu\left(p_{z}, p, \bar{s}_{p}, d\right)=\bar{s}_{p}-1 / p+1 / 2$ is strictly less than $z=3 / 2$ if, and only if, $\bar{s}_{p}<1+1 / p$. In this case, also Formula (5) in Theorem 2.1 applies which proves the upper bound on $\bar{\alpha}_{p}$ in case 1.). Hence, the proof is complete.

Let us add some remarks also for this example.
Remark 3.7. There exist statements similar to Proposition 3.5 also for $p \geq 2$; see, e.g., Ebmeyer [14] for details. However, in this case the analogue of (12) does not provide additional information; cf. Remark 2.4(iv). That is, using Theorem 2.1 not much can be said except that $\bar{\alpha}_{p}(S(\Omega)) \geq \bar{s}_{p}(S(\Omega))$ might be unbounded. Anyway, again this agrees well with results due to Dahlke [5], who showed that for $p=d=2$ and smooth right-hand sides we indeed have $\bar{\alpha}_{2}(S(\Omega))=\infty>\bar{s}_{2}(S(\Omega))$; see also Remark 3.3 above.
Remark 3.8. Theorem 3.6 shows that on polyhedral Lipschitz domains the maximal $L_{p}(\Omega)$ Sobolev smoothness $\bar{s}_{p}$ is at least $3 / 2$. In [30, Theorem 2'] Savaré proved that this remains true on general Lipschitz domains under the weaker condition that $f \in W_{p^{\prime}}^{-1 / 2}(\Omega)$. Moreover, in [30, Remark 4.3] he even claims optimality. However, if we stick to the stronger assumptions that $\Omega$ is polyhedral Lipschitz and $f \in L_{p^{\prime}}(\Omega)$, we may use positive Besov regularity results w.r.t. the scale $(*)$ in order to conclude a better lower bound. Indeed, combining Proposition 3.5 with Remark 2.4(v) shows that

$$
\bar{s}_{p} \geq \widetilde{s}_{p} \geq \alpha \cdot \frac{z+d\left(1 / p-1 / p_{z}\right)}{\alpha+d\left(1 / p-1 / p_{z}\right)}=\left(1+\frac{1}{p}\right) \frac{\alpha}{\alpha+1 / p-1 / 2}=: \widehat{s}_{p}(\alpha) \quad \text { for all } \quad \alpha \leq \bar{\alpha}_{p}
$$

Note that this lower bound is strictly monotonically increasing in $\alpha$, where

$$
\frac{3}{2}=\widehat{s}_{p}(3 / 2)<\widehat{s}_{p}(\alpha)<1+\frac{1}{p}, \quad \alpha>\frac{3}{2} .
$$

Results of Dahlke et al. [8, Theorem 4.20] imply that on bounded polygonal domains $\Omega \subseteq \mathbb{R}^{2}$,

$$
\begin{equation*}
S(\Omega):=\left\{u \in W_{p, 0}^{1}(\Omega) \mid \Delta_{p} u \in L_{\infty}(\Omega)\right\} \subseteq B_{\tau, \tau}^{\alpha}(\Omega), \quad \frac{1}{\tau}=\frac{\alpha}{2}+\frac{1}{p}, \quad \text { for all } \quad 0<\alpha<2 \tag{13}
\end{equation*}
$$

such that in this case

$$
\bar{s}_{p}(S(\Omega)) \geq 2 \frac{1+1 / p}{1+1 / p+1 / 2}, \quad 1<p \leq 2
$$

Furthermore, recent results indicate that we may replace $L_{\infty}(\Omega)$ by $L_{p^{\prime}}(\Omega)$ in (13).

### 3.3 The inhomogeneous stationary Stokes problem

Our third and final example is the inhomogeneous stationary Stokes system

$$
\left.\begin{array}{rlc}
-\Delta u+\nabla \pi=f & \text { in } & \Omega,  \tag{14}\\
\operatorname{div}(u)=g & \text { in } & \Omega, \\
u_{\partial \Omega}=h & \text { on } & \partial \Omega,
\end{array}\right\}
$$

where $\Omega \subseteq \mathbb{R}^{d}$ is again a bounded Lipschitz domain $(d \geq 2)$ and $f, g$, and $h$ are given functions (or distributions) on $\Omega$ and $\partial \Omega$, respectively, such that the compatibility condition

$$
\begin{equation*}
\int_{\partial \Omega} h(y) \cdot \eta(y) \mathrm{d} y=\int_{\Omega} g(x) \mathrm{d} x \tag{15}
\end{equation*}
$$

is satisfied; here, $\eta$ denotes the outward unit normal vector to $\partial \Omega$.
For this problem, Mitrea and Wright [28] showed that a suitably modified regularity shift holds in a range of parameters $\mathcal{R}_{d, \varepsilon} \subseteq \mathbb{R} \times(0, \infty]$ similar to the one established by Jerison and Kenig [23] for the classical Poisson problem; see [28, page 178] for a precise definition of $\mathcal{R}_{d, \varepsilon}$. Without going into details, this range depends on a "roughness parameter" $\varepsilon=\varepsilon(\Omega) \in(0,1]$ which measures the Lipschitz nature of $\Omega$. However, for sufficiently smooth domains, e.g., when $\partial \Omega \in \mathcal{C}^{1}$, we may take $\varepsilon=1$.

Proposition 3.9 (Mitrea and Wright [28, Theorem 1.5/10.15]). For $d \geq 2$ let $\Omega \subseteq \mathbb{R}^{d}$ be a bounded Lipschitz domain. Moreover, let $A \in\{B, F\}$, as well as $(d-1) / d<p \leq \infty$, $0<q \leq \infty$, and $(d-1) \max \{1 / p-1,0\}<s<1$ with $(s, p) \in \mathcal{R}_{d, \varepsilon(\Omega)}$, where $\max \{p, q\}<\infty$ if $A=F$. Then for

$$
f \in A_{p, q}^{s+1 / p-2}(\Omega)^{d}, \quad g \in A_{p, q}^{s+1 / p-1}(\Omega), \quad \text { and } \quad h \in \begin{cases}B_{p, q}^{s}(\partial \Omega)^{d} & \text { if } A=B \\ F_{p, p}^{s}(\partial \Omega)^{d} & \text { if } A=F\end{cases}
$$

there exists a solution $(u, \pi) \in A_{p, q}^{s+1 / p}(\Omega)^{d} \times A_{p, q}^{s+1 / p-1}(\Omega)$ to the inhomogeneous stationary Stokes system (14), (15). Moreover, it is unique modulo the addition of locally constant functions in $\Omega$ to the pressure $\pi$.

This statement can be used to conclude the subsequent regularity assertion which provides all necessary information for the application of Theorem 2.1 to the Stokes problem.

Lemma 3.10. For $d \geq 2$ let $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded Lipschitz domain with roughness parameter $\varepsilon(\Omega) \in(0,1]$. Further, let $0<s<1$, as well as $\sigma:=\min _{j \in\{1,2,3\}} \sigma_{j} \geq 0$ and

$$
f \in H^{s-3 / 2+\sigma_{1}}(\Omega)^{d}, \quad g \in H^{s-1 / 2+\sigma_{2}}(\Omega), \quad \text { and } \quad h \in H^{s+\sigma_{3}}(\partial \Omega)^{d}
$$

Then solutions $(u, \pi)$ to (14), (15) exist and satisfy $(u, \pi) \in H_{p}^{s+1 / p}(\Omega)^{d} \times H_{p}^{s+1 / p-1}(\Omega)$ for all $p \in[2, \infty)$ with

$$
\begin{equation*}
\frac{1}{2}-\min \left\{\frac{\varepsilon(\Omega)}{2}, \frac{\sigma}{d-1}\right\} \leq \frac{1}{p} \leq \frac{1}{2} \tag{16}
\end{equation*}
$$

Proof. Due to simple embeddings we may w.l.o.g. assume that $0 \leq \sigma<1 / 2$; see Proposition A.3(iv). Further let $s \in \mathbb{R}$ and $p \in[2, \infty)$. Then, according to Definition A. 2 and Proposition A.3, there holds

$$
H^{s-3 / 2+\sigma_{1}}(\Omega) \hookrightarrow F_{2,2}^{s-3 / 2+\sigma}(\Omega) \hookrightarrow F_{p, 2}^{s_{1}}(\Omega) \hookrightarrow F_{p, 2}^{s+1 / p-2}(\Omega)
$$

provided that

$$
s_{1}:=s+\frac{1}{p}-2+\sigma+(d-1)\left(\frac{1}{p}-\frac{1}{2}\right) \geq s+\frac{1}{p}-2 .
$$

Note that this inequality is satisfied if $p$ is chosen such that

$$
\begin{equation*}
\frac{1}{2}-\frac{\sigma}{d-1} \leq \frac{1}{p} \tag{17}
\end{equation*}
$$

Moreover, similar calculations show that the same condition (17) implies the embeddings $H^{s-1 / 2+\sigma_{2}}(\Omega) \hookrightarrow F_{p, 2}^{s+1 / p-1}(\Omega)$ and $H^{s+\sigma_{3}}(\partial \Omega) \hookrightarrow F_{p, p}^{s}(\partial \Omega)$. Hence, our assumptions on the data give

$$
f \in F_{p, 2}^{s+1 / p-2}(\Omega)^{d}, \quad g \in F_{p, 2}^{s+1 / p-1}(\Omega), \quad \text { and } \quad h \in F_{p, p}^{s}(\partial \Omega)^{d}
$$

with $0<s<1$ and each $p \in[2, \infty)$ with (17). Furthermore, it can be checked easily that $(s, p) \in \mathcal{R}_{d, \varepsilon(\Omega)}$ whenever $0<s<1$ and $p \in[2, \infty)$ with

$$
\frac{1}{2}-\frac{\varepsilon(\Omega)}{2} \leq \frac{1}{p}
$$

Thus, the claim follows from Proposition 3.9 applied for $A:=F, q:=2$, as well as $0<s<1$ and $p \in[2, \infty)$ restricted by (16), and Definition A.2.

Theorem 3.11. For $d \geq 2$ let $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded Lipschitz domain with roughness parameter $\varepsilon=\varepsilon(\Omega) \in(0,1]$. Let $S_{u}(\Omega)$ and $S_{\pi}(\Omega)$ denote the sets of solutions $(u, \pi)$ to the inhomogeneous stationary Stokes problem (14), (15) with

$$
f \in H^{-1 / 2+\sigma_{1}}(\Omega)^{d}, \quad g \in H^{1 / 2+\sigma_{2}}(\Omega), \quad \text { and } \quad h \in H^{1+\sigma_{3}}(\partial \Omega)^{d},
$$

where

$$
\sigma:=\min _{j \in\{1,2,3\}} \sigma_{j}>0 .
$$

Moreover, let $m:=\min \{(d-1) \varepsilon / 2, \sigma\}$. Then for the regularity indices $\bar{s}_{2}:=\bar{s}_{2}\left(S_{u}(\Omega)\right)$ and $\bar{\alpha}_{2}:=\bar{\alpha}_{2}\left(S_{u}(\Omega)\right)$ of (each component of) the velocity $u$ one of the following cases applies:
1.) $3 / 2 \leq \bar{s}_{2}<3 / 2+m$ and

$$
\bar{s}_{2} \leq \bar{\alpha}_{2} \leq \bar{s}_{2} \cdot \frac{d}{d-1} \cdot \frac{m}{3 / 2+m-\bar{s}_{2}}
$$

2.) $3 / 2+m \leq \bar{s}_{2} \leq \bar{\alpha}_{2}$.

For the regularity of the pressure $\pi$ an analogous statement holds with $3 / 2$ replaced by $1 / 2$.
Proof. Let us only consider the assertions on $S_{u}(\Omega)$; the results for $S_{\pi}(\Omega)$ can be derived exactly in the same way. Due to Proposition A.3(iv) and Lemma 3.10 applied for $p=2$ we know that

$$
S_{u}(\Omega) \subseteq H_{p}^{s+1 / 2}(\Omega)^{d}=F_{p, 2}^{s+1 / 2}(\Omega)^{d} \quad \text { for all } \quad s<1
$$

Therefore, by Remark 2.4(i) we have $3 / 2 \leq \bar{s}_{2} \leq \bar{\alpha}_{2}$.
Since $m>0$, it remains to show that if $\bar{s}_{2}<3 / 2+m$, then the stated upper bound on $\bar{\alpha}_{2}$ holds true. To this end, let us define

$$
\bar{\delta}:=\min \left\{1, \frac{1}{2}\left(\frac{3}{2}+m-\bar{s}_{2}\right)\right\} .
$$

Then $3 / 2 \leq \bar{s}_{2}<3 / 2+m$ particularly implies that $0<\bar{\delta}<m \leq(d-1) / 2$. For each arbitrarily fixed $\delta \in(0, \bar{\delta})$ we can now choose $p_{z}=p_{z}(\delta) \in(2, \infty)$ with

$$
(d-1)\left(\frac{1}{2}-\frac{1}{p_{z}}\right)=m-\delta
$$

Then the definition of $m$ implies that

$$
0<\frac{1}{2}-\frac{1}{p_{z}}<\min \left\{\frac{\varepsilon(\Omega)}{2}, \frac{\sigma}{d-1}\right\}
$$

and hence $p_{z}$ satisfies (16). Thus, Lemma 3.10 ensures that $S_{u}(\Omega) \subseteq H_{p_{z}}^{s}(\Omega)=F_{p_{z}, 2}^{s}(\Omega)$ for all $s<z:=1+1 / p_{z}$. According to Remark 2.2, this allows to apply Theorem 2.1, where

$$
\mu=\bar{s}_{2}-d\left(\frac{1}{2}-\frac{1}{p_{z}}\right)=z+\bar{s}_{2}-\frac{3}{2}-m+\delta<z-\frac{1}{2}\left(\frac{3}{2}+m-\bar{s}_{2}\right)<z
$$

Therefore, the bound (5) applies which shows that

$$
\bar{\alpha}_{2} \leq \bar{s}_{2} \cdot \frac{\bar{s}_{2}-\mu}{z-\mu}=\bar{s}_{2} \cdot \frac{d\left(1 / 2-1 / p_{z}\right)}{3 / 2+m-\delta-\bar{s}_{2}}=\bar{s}_{2} \cdot \frac{d}{d-1} \cdot \frac{m-\delta}{3 / 2+m-\delta-\bar{s}_{2}} .
$$

Since the latter inequality holds for arbitrary small $\delta>0$, this completes the proof.

Let us conclude also this section with some final remarks:
Remark 3.12. Assume for simplicity that $\sigma=\sigma_{1}$ is chosen small enough such that $m=\sigma$. Then case 2.) in Theorem 3.11 can be interpreted as a shift $H^{-1 / 2+\sigma} \ni f \mapsto u \in H^{3 / 2+\sigma}$ of full order (two) within the Sobolev scale. However, as we have seen in Section 3.1, already for the classical Poisson problem this shift might fail even on $\mathcal{C}^{1}$ domains. Although we do not know about an explicit example, it is very likely that the same is true for the Stokes problem. Then case 1.) applies and we have a non-trivial upper bound $\bar{\alpha}_{2} \leq b$ on the Besov smoothness w.r.t. the scale $(*)$ with $p=2$. Moreover note that this $b=b\left(\bar{s}_{2}\right)$ is monotonically increasing in $\bar{s}_{2}$, where

$$
\frac{3}{2} \frac{d}{d-1}=b(3 / 2) \leq b\left(\bar{s}_{2}\right)<b(3 / 2+m)=\infty, \quad \bar{s}_{2} \in[3 / 2,3 / 2+m)
$$

Recently Eckhardt et al. [16, Theorem 3.3] addressed the question of Besov regularity for dimensions $d \geq 3$ under the additional conditions that the boundary of $\Omega$ is connected and $g=0$. Rewritten in our notation they were able to show that for $\sigma_{1}=1 / 2$ and $\sigma_{3}=0$ we have for $d \geq 4$

$$
\bar{\alpha}_{2}\left(S_{u}(\Omega)\right) \geq \frac{3}{2} \frac{d}{d-1} \quad \text { and } \quad \bar{\alpha}_{2}\left(S_{\pi}(\Omega)\right) \geq \frac{1}{2} \frac{d}{d-1} .
$$

## A Appendix: Basics from function space theory

In this supplementary section we collect the main definitions and assertions concerning function spaces on domains which are needed throughout the paper. Here 'domain' always means 'non-empty, connected, open set'. Special attention is paid to bounded Lipschitz domains $\Omega \subseteq \mathbb{R}^{d}, d \in \mathbb{N}$, as defined, e.g., in Triebel [32, Section 1.11.4].

## A. 1 Besov and Triebel-Lizorkin spaces

In accordance with Triebel [31] we use the Fourier analytic approach towards Besov and Triebel-Lizorkin spaces on $\mathbb{R}^{d}$ and define the corresponding spaces on domains by restriction.

Let $d \in \mathbb{N}$. By $\mathcal{S}\left(\mathbb{R}^{d}\right)$ we denote the Schwartz space of all complex-valued rapidly decreasing $\mathcal{C}^{\infty}$ functions on $\mathbb{R}^{d}$ and $\mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$ denotes its dual space of tempered distributions. Moreover, for domains $\Omega \subseteq \mathbb{R}^{d}$ we let $\mathcal{D}(\Omega):=\mathcal{C}_{0}^{\infty}(\Omega)$ denote the collection of all complexvalued $\mathcal{C}^{\infty}$ functions in $\mathbb{R}^{d}$ with compact support in $\Omega$ and denote by $\mathcal{D}^{\prime}(\Omega)$ its dual space of distributions on $\Omega$. As usual, we say two functionals $f$ and $g$ equal each other in $\mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$ or $\mathcal{D}^{\prime}(\Omega)$ if

$$
f(\varphi)=g(\varphi) \quad \text { for all } \varphi \text { from } \mathcal{S}\left(\mathbb{R}^{d}\right) \text { or } \mathcal{D}(\Omega), \text { respectively. }
$$

For $g \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$ we denote by $g_{\left.\right|_{\Omega}}$ the restriction of $g$ to $\Omega$ which means that

$$
g_{\left.\right|_{\Omega}} \in \mathcal{D}^{\prime}(\Omega) \quad \text { and } \quad\left(g_{\left.\right|_{\Omega}}\right)(\varphi):=g(\varphi) \quad \text { for all } \quad \varphi \in \mathcal{D}(\Omega)
$$

Note that this is meaningful since $\mathcal{D}(\Omega) \subseteq \mathcal{D}\left(\mathbb{R}^{d}\right) \subseteq \mathcal{S}\left(\mathbb{R}^{d}\right)$.
In addition, let $\mathcal{F}$ and $\mathcal{F}^{-1}$ denote the (extension of the) Fourier transform, respectively its inverse, on $\mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$. Fix an arbitrary $\phi_{0} \in \mathcal{S}\left(\mathbb{R}^{d}\right)$ such that

$$
\phi_{0}(x)=1 \quad \text { if } \quad|x|_{2} \leq 1 \quad \text { and } \quad \phi_{0}(x)=0 \quad \text { if } \quad|x|_{2} \geq \frac{3}{2} .
$$

Then the collection $\Phi:=\left(\phi_{k}\right)_{k \in \mathbb{N}_{0}}$, with

$$
\phi_{k}(x):=\phi_{0}\left(2^{-k} x\right)-\phi_{0}\left(2^{-k+1} x\right), \quad x \in \mathbb{R}^{d}, \quad k \in \mathbb{N},
$$

defines a smooth dyadic resolution of unity and we have

$$
f=\sum_{k=0}^{\infty} \mathcal{F}^{-1}\left[\phi_{k} \mathcal{F} f\right] \quad\left(\text { convergence in } \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)\right)
$$

for all $f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$. Due to the celebrated Paley-Wiener-Schwartz-Theorem, the building blocks $\mathcal{F}^{-1}\left[\phi_{k} \mathcal{F} f\right], k \in \mathbb{N}_{0}$, are actually entire analytic functions; see, for instance, Triebel [31, Section 1.2.1]. As usual, for $0<q<\infty, \ell_{q}\left(\mathbb{N}_{0}\right)$ is the space of $q$-summable scalar-valued sequences over $\mathbb{N}_{0}$ (bounded sequences, if $q=\infty$ ).

Definition A.1. For $d \in \mathbb{N}$ choose $\Phi$ as above and let $\Omega \subsetneq \mathbb{R}^{d}$ denote an arbitrary domain. Moreover, let $s \in \mathbb{R}$ and $0<p, q \leq \infty$.
(i) The set $B_{p, q}^{s}\left(\mathbb{R}^{d}\right):=\left\{f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)\left|\left\|f \mid B_{p, q}^{s}\left(\mathbb{R}^{d}\right)\right\|<\infty\right\}\right.$, quasi-normed by

$$
\left\|f\left|B_{p, q}^{s}\left(\mathbb{R}^{d}\right)\|:=\|\left(2^{k s}\left\|\mathcal{F}^{-1}\left[\phi_{k} \mathcal{F} f\right](\cdot) \mid L_{p}\left(\mathbb{R}^{d}\right)\right\|\right)_{k \in \mathbb{N}_{0}}\right| \ell_{q}\left(\mathbb{N}_{0}\right)\right\|,
$$

is called Besov space.
(ii) If $p<\infty$, then the set $F_{p, q}^{s}\left(\mathbb{R}^{d}\right):=\left\{f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)\left|\left\|f \mid F_{p, q}^{s}\left(\mathbb{R}^{d}\right)\right\|<\infty\right\}\right.$, quasi-normed by

$$
\left\|f \left|F _ { p , q } ^ { s } ( \mathbb { R } ^ { d } ) \| : = \| \left\|( 2 ^ { k s } | \mathcal { F } ^ { - 1 } [ \phi _ { k } \mathcal { F } f ] ( \cdot ) | ) _ { k \in \mathbb { N } _ { 0 } } \left|\ell_{q}\left(\mathbb{N}_{0}\right)\left\|\mid L_{p}\left(\mathbb{R}^{d}\right)\right\|,\right.\right.\right.\right.
$$

is called Triebel-Lizorkin space.
(iii) If $A \in\{B, F\}$ with $p<\infty$ for $A=F$, then the set

$$
A_{p, q}^{s}(\Omega):=\left\{f \in \mathcal{D}^{\prime}(\Omega) \mid \text { there exists } g \in A_{p, q}^{s}\left(\mathbb{R}^{d}\right) \text { with } g_{\left.\right|_{\Omega}}=f \text { in } \mathcal{D}^{\prime}(\Omega)\right\}
$$

quasi-normed by

$$
\left\|f\left|A_{p, q}^{s}(\Omega)\left\|:=\inf _{\substack{g \in A_{p, q}^{s}\left(\mathbb{R}^{d}\right) \\ g_{\left.\right|_{\Omega}}=f \text { in } \mathcal{D}^{\prime}(\Omega)}}\right\| g\right| A_{p, q}^{s}\left(\mathbb{R}^{d}\right)\right\|,
$$

is called Besov resp. Triebel-Lizorkin space on $\Omega$.
Standard proofs show that the spaces introduced above are quasi-Banach spaces (Banach iff $\min \{p, q\} \geq 1$ and Hilbert iff $p=q=2$ ) and that different $\Phi$ provide equivalent quasinorms, see, e.g., Triebel [31, Section 2.3.2]. Furthermore, these scales of spaces cover a variety of classical function spaces - such as, e.g., Lebesgue, Sobolev(-Slobodeckij), Bessel potential, Lipschitz, Hölder(-Zygmund), or Hardy spaces - as special cases. Besides our Fourier analytic definition, there is a big variety of other descriptions of these spaces which are equivalent at least for large ranges of parameters. To give an example, we note that at least for

$$
s>\sigma_{p}:=d \max \left\{\frac{1}{p}-1,0\right\}
$$

the spaces $A_{p, q}^{s}\left(\mathbb{R}^{d}\right)$ (and also $A_{p, q}^{s}(\Omega)$ for bounded Lipschitz domains $\Omega \subseteq \mathbb{R}^{d}$ ) exclusively contain regular distributions, i.e., functions, which makes it possible to characterize them as subspaces of some Lebesgue space by means of iterated differences. For details we refer to Triebel [32, Section 1.11.9].

## A. 2 Sobolev spaces

We follow the usual approach and define the subsequent Sobolev-type spaces based on Besov and Triebel-Lizorkin spaces.

Definition A.2. For $d \in \mathbb{N}$ let $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded Lipschitz domain. Then we set

$$
\begin{array}{rrl}
W_{p}^{m}(\Omega):=F_{p, 2}^{m}(\Omega), & m \in \mathbb{N}_{0}, 1<p<\infty, & \text { (Sobolev) } \\
W_{p}^{s}(\Omega):=F_{p, p}^{s}(\Omega)=B_{p, p}^{s}(\Omega), & 0<s \notin \mathbb{N}, 1 \leq p<\infty, & \text { (Sobolev-Slobodeckij) } \\
W_{p}^{s}(\Omega):=\left[W_{p^{\prime}, 0}^{-s}(\Omega)\right]^{\prime}, & s<0,1<p<\infty, \\
H_{p}^{s}(\Omega):=F_{p, 2}^{s}(\Omega), & s \in \mathbb{R}, 1<p<\infty, & \text { (Bessel potential) } \\
H^{s}(\Omega):=H_{2}^{s}(\Omega)=F_{2,2}^{s}(\Omega)=B_{2,2}^{s}(\Omega), & s \in \mathbb{R}, & \text { (Sobolev-Hilbert) }
\end{array}
$$

where for $1<p<\infty$ the index $p^{\prime}$ is given by $1 / p+1 / p^{\prime}=1$ and $W_{p, 0}^{s}(\Omega)$ denotes the closure of $\mathcal{C}_{0}^{\infty}(\Omega)$ w.r.t. the norm $\left\|\cdot \mid W_{p}^{s}(\Omega)\right\|$ if $s>0$.

It is worth noting that these definitions are equivalent with the common definitions of Sobolev(-Slobodeckij) and Bessel potential spaces. In particular, for $s=m \in \mathbb{N}_{0}$ we have

$$
W_{p}^{m}(\Omega)=\left\{f \in L_{p}(\Omega)\left|\left\|f \mid W_{p}^{m}(\Omega)\right\|:=\left[\sum_{|\alpha|_{1} \leq m}\left\|D^{\alpha} f \mid L_{p}(\Omega)\right\|^{p}\right]^{1 / p}<\infty\right\}\right.
$$

see Triebel [32, Theorem 1.122], while $W_{p}^{s}(\Omega)=B_{p, p}^{s}(\Omega)$ for $0<s \notin \mathbb{N}$ coincides with the classical definition of Sobolev-Slobodeckij spaces as real interpolation space of $L_{p}(\Omega)$ with $W_{p}^{m}(\Omega)$ for some $s<m \in \mathbb{N}$ and suitable parameters; see, e.g., DeVore [13, Section 4.6].

## A. 3 Embeddings

The scales of Besov and Triebel-Lizorkin spaces $A_{p, q}^{s}(\Omega)$ on bounded Lipschitz domains satisfy various embeddings. Let us mention a few of them:

Proposition A.3. For $d \in \mathbb{N}$ let $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded Lipschitz domain. Further assume $s, s_{0}, s_{1} \in \mathbb{R}$ and let $0<p, p_{0}, p_{1}, q, q_{0}, q_{1} \leq \infty$.
(i) Assume additionally that $p<\infty$. Then

$$
B_{p, q_{0}}^{s}(\Omega) \hookrightarrow F_{p, q}^{s}(\Omega) \hookrightarrow B_{p, q_{1}}^{s}(\Omega)
$$

holds if, and only if, we have $q_{0} \leq \min \{p, q\} \leq \max \{p, q\} \leq q_{1}$.
(ii) If additionally $p_{0}<p_{1}<\infty$ and $s_{0}-d / p_{0}=s_{1}-d / p_{1}$, then

$$
F_{p_{0}, q_{0}}^{s_{0}}(\Omega) \hookrightarrow F_{p_{1}, q_{1}}^{s_{1}}(\Omega)
$$

(iii) If additionally $A \in\{B, F\}$ (and $p<\infty$ if $A=F$ ), as well as $q_{0} \leq q_{1}$, then

$$
A_{p, q_{0}}^{s}(\Omega) \hookrightarrow A_{p, q_{1}}^{s}(\Omega)
$$

(iv) If additionally $X, Y \in\{B, F\}$ and

$$
s_{0}-s_{1}>d \max \left\{\frac{1}{p_{0}}-\frac{1}{p_{1}}, 0\right\},
$$

then

$$
X_{p_{0}, q_{0}}^{s_{0}}(\Omega) \hookrightarrow Y_{p_{1}, q_{1}}^{s_{1}}(\Omega)
$$

(with finite integrability parameter for $F$-spaces).
(v) Assume additionally that $p_{0}<p<p_{1}$ and

$$
s_{0}-\frac{d}{p_{0}}=s-\frac{d}{p}=s_{1}-\frac{d}{p_{1}} .
$$

Then

$$
B_{p_{0}, q_{0}}^{s_{0}}(\Omega) \hookrightarrow F_{p, q}^{s}(\Omega) \hookrightarrow B_{p_{1}, q_{1}}^{s_{1}}(\Omega)
$$

holds if, and only if, we have $q_{0} \leq p \leq q_{1}$.
Proof. For (i), (ii), and (v) see, e.g., Triebel [32, page 60] and the references therein. For (iii) and (iv) additionally consult Triebel [31, Proposition 2 in Section 2.3.2], as well as [33, Theorem 4.33 and Remark 4.34].

Note that Proposition A.3(iv) particularly implies that for $A \in\{B, F\}$ we have

$$
A_{p_{0}, q}^{s_{0}}(\Omega) \hookrightarrow W_{p_{1}}^{s_{1}}(\Omega) \quad \text { if } \quad s_{0}>s_{1} \geq 0, \text { as well as } 1<p_{1} \leq p_{0} \leq \infty, \text { and } 0<q \leq \infty
$$

with $p_{0}<\infty$ if $A=F$, since $W_{p}^{s_{1}}(\Omega)$ can be identified with $F_{p, q}^{s_{1}}(\Omega)$ for $q \in\{2, p\}$.

## A. 4 Complex interpolation

For some open set $\Omega$ let $X(\Omega)$ and $Y(\Omega)$ denote quasi-normed spaces of complex-valued functions or distributions on $\Omega$. Then, under certain conditions, the application of the so-called (extended) complex interpolation method with parameter $\theta \in(0,1)$ defines another quasi-normed space of functions on $\Omega$. Besides other useful properties this space, usually denoted by $[X(\Omega), Y(\Omega)]_{\theta}$, satisfies

$$
X(\Omega) \cap Y(\Omega) \hookrightarrow[X(\Omega), Y(\Omega)]_{\theta} \hookrightarrow X(\Omega)+Y(\Omega) .
$$

For details we refer to Bergh, Löfström [2] and Kalton, Mayboroda, Mitrea [24].
It turns out that the scales of Besov and Triebel-Lizorkin spaces $A_{p, q}^{s}(\Omega)$ on bounded Lipschitz domains behave well w.r.t. this method:
Proposition A. 4 (Kalton et al. [24, Theorem 9.4]). For $d \in \mathbb{N}$ let $\Omega \subseteq \mathbb{R}^{d}$ denote a bounded Lipschitz domain and assume $\theta \in(0,1)$. Moreover, let $A \in\{B, F\}$, as well as $s, s_{0}, s_{1} \in \mathbb{R}$, and $0<p, p_{0}, p_{1}, q, q_{0}, q_{1} \leq \infty$ (with $p_{0}, p_{1}<\infty$ for $A=F$ ), and $\min \left\{q_{0}, q_{1}\right\}<\infty$. Then

$$
s=(1-\theta) s_{0}+\theta s_{1}, \quad \frac{1}{p}=\frac{1-\theta}{p_{0}}+\frac{\theta}{p_{1}}, \quad \text { and } \quad \frac{1}{q}=\frac{1-\theta}{q_{0}}+\frac{\theta}{q_{1}}
$$

implies

$$
\left[A_{p_{0}, q_{0}}^{s_{0}}(\Omega), A_{p_{1}, q_{1}}^{s_{1}}(\Omega)\right]_{\theta}=A_{p, q}^{s}(\Omega)
$$

in the sense of equivalent quasi-norms.

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